

Table 5.1 from (2002TI10): Energy levels of  ${}^5\text{He}$ , extended  $R$ -matrix prescription <sup>a</sup>

$E_x$ (MeV)	$J^\pi; T$	$\Gamma_{\text{cm}}$ <sup>b</sup> (MeV)	$\Gamma_n$ (MeV)	$\Gamma_d$ (MeV)	$\Gamma_{n^*}$ <sup>c</sup> (MeV)	Decay	Reactions (used in analysis)
g.s. <sup>d</sup>	$\frac{3}{2}^-; \frac{1}{2}$	0.648	0.578	8.80 <sup>e</sup>	66.0 <sup>e</sup>	n, $\alpha$	5, 8, 13, 23, 24, 25
1.27	$\frac{1}{2}^-; \frac{1}{2}$	5.57	3.18	38.0 <sup>e</sup>	1.27 <sup>e</sup>	n, $\alpha$	5, 8, 21, 24, 25
16.84	$\frac{3}{2}^+; \frac{1}{2}$	0.0745	0.040	0.025 <sup>f</sup>		$\gamma$ , n, d, t, $\alpha$	2, 3, 7, 8, 10, 13, 14, 23, 24, 25
19.14	$\frac{5}{2}^+; \frac{1}{2}$	3.56	0.003	1.62 <sup>g</sup>		n, d, t, $\alpha$	4, 10, 14, 23
19.26	$\frac{3}{2}^+; \frac{1}{2}$	3.96	0.014	1.83 <sup>g</sup>		n, d, t, $\alpha$	4, 10, 14, 23
19.31	$\frac{7}{2}^+; \frac{1}{2}$	3.02	0.045	1.89 <sup>g</sup>		n, d, t, $\alpha$	4, 10, 14
19.96	$\frac{3}{2}^-; \frac{1}{2}$	1.92	0.003	0.325 <sup>h</sup>	0.862	n, p, d, t, $\alpha$	3, 17, 24, 25
21.25	$\frac{3}{2}^+; \frac{1}{2}$	4.61	0.098	2.38 <sup>i</sup>		n, d, t, $\alpha$	21
21.39	$\frac{5}{2}^+; \frac{1}{2}$	3.95	0.091	2.12 <sup>l</sup>		n, d, t, $\alpha$	21
21.64	$\frac{1}{2}^+; \frac{1}{2}$	4.03	0.050	0.878 <sup>j</sup>	0.726	n, p, d, t, $\alpha$	21
23.97	$\frac{7}{2}^+; \frac{1}{2}$	5.44	0.053	2.85 <sup>g</sup>		n, d, t, $\alpha$	
24.06	$\frac{5}{2}^-; \frac{1}{2}$	5.23	0.013	2.18 <sup>k</sup>		n, d, t, $\alpha$	
$(35.7 \pm 0.4)$ <sup>l</sup>		$\approx 2$ <sup>l</sup>					21, 25

<sup>a</sup> This prescription, based on the complex poles and residues of the  $S$ -matrix, is the recommended one (see Introduction). The channel radii are:  $a_n = 3.0$  fm,  $a_d = 5.1$  fm. The uncertainties in the widths and positions of the first three levels are less than 1%. Above 19 MeV excitation energy, they increase rapidly, varying from about 5% up to as much as 50% for the broad higher levels. Except where noted, all parameters in the table are newly adopted in this evaluation.

<sup>b</sup> The fact that the sum of the partial widths is unequal to the total width in the extended  $R$ -matrix prescription is characteristic of non-Breit-Wigner resonances as was discussed in the appendix of (1992TI02).

<sup>c</sup> The  $n^*$  designation indicates  $n + \alpha^*$  where the first excited state of the  $\alpha$  particle was included as a way to approximate the effects of three-body breakup on the two-body channels.

<sup>d</sup> Situated 798 keV above the  $n + \alpha$  threshold. This value is in excellent agreement with early measurements reported by (1963SM03;  $790 \pm 30$  keV) and by (1960YO06;  $800 \pm 100$  keV).

<sup>e</sup> These large partial widths in closed channels have no meaning as decay widths, but rather as asymptotic normalization constants.

<sup>f</sup> Entirely  ${}^4S(d)$ .

<sup>g</sup> Primarily  ${}^4D(d)$ .

<sup>h</sup> Primarily  ${}^2P(d)$ .

<sup>i</sup> Primarily  ${}^2D(d)$ .

<sup>j</sup> Primarily  ${}^2S(d)$ .

<sup>k</sup> Primarily  ${}^4P(d)$ .

<sup>l</sup> Retained from the previous evaluation (1988AJ01).