Energy Levels of Light Nuclei

$A = 8$

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Abstract: An evaluation of $A = 5$–10 was published in *Nuclear Physics A227* (1974), p. 1. This version of $A = 8$ differs from the published version in that we have corrected some errors discovered after the article went to press. Figures and introductory tables have been omitted from this manuscript. Reference key numbers have been changed to the TUNL/NNDC format.

(References closed December 31, 1973)

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$^8\text{He}$

(Figs. 11 and 14)

GENERAL: (See also (1966LA04).)


Mass of $^8\text{He}$: The atomic mass excess of $^8\text{He}$ derived from the $Q$ of the $^{26}\text{Mg}(\alpha, ^8\text{He})^{22}\text{Mg}$ reaction is $31.65 \pm 0.12$ MeV. See also (1968BA48). $^8\text{He}$ is then stable to decay into $^6\text{He} + 2n$ by 2.1 MeV (1966CE01). See also (1966BA38, 1968CE1A, 1972CE1A).

1. $^8\text{He}(\beta^-)^8\text{Li}$

$Q_m = 10.70$

The half-life is $122 \pm 2$ msec. The decay takes place $88\%$ to $^8\text{Li}^*(0.98)$ [log $ft = 4.20$; B. Zimmerman, private communication] and $12 \pm 1\%$ via $^8\text{Li}$ states decaying by neutron emission (1965PO06). See also (1966NE07) and (1973HA49; theor.).

2. $^{11}\text{B}(\gamma, 3p)^8\text{He}$

$Q_m = -44.85$

See (1966NE07).

3. $^{26}\text{Mg}(\alpha, ^{22}\text{Mg})^8\text{He}$

$Q_m = -45.05$

This reaction has been studied at $E_{\alpha} = 80$ MeV (1966CE01).

$^8\text{Li}$

(Figs. 11 and 14)

GENERAL: (See also (1966LA04).)


Special levels: (1966BA26, 1970FR1C).

Table 8.1: Energy levels of $^8$Li

<table>
<thead>
<tr>
<th>$E_x$ (MeV ± keV)</th>
<th>$J^\pi; T$</th>
<th>$\tau_m$ or $\Gamma_{c.m.}$ (keV)</th>
<th>Decay</th>
<th>Reactions</th>
</tr>
</thead>
<tbody>
<tr>
<td>g.s.</td>
<td>$2^+; 1$</td>
<td>$\tau_{1/2} = 842 \pm 6$ msec</td>
<td>$\beta^-$</td>
<td>1, 2, 3, 10, 11, 14, 17, 18, 19, 20, 21, 23, 24, 25, 26, 27</td>
</tr>
<tr>
<td>0.9808 ± 0.1</td>
<td>$1^+; 1$</td>
<td>$\tau_m = 12 \pm 4$ fsec</td>
<td>$\gamma$</td>
<td>2, 3, 10, 13, 14, 15, 17, 18, 19, 21, 24, 25, 26</td>
</tr>
<tr>
<td>2.261 ± 2</td>
<td>$3^+; 1$</td>
<td>$\Gamma = 31 \pm 5$</td>
<td>$\gamma, n$</td>
<td>3, 4, 10, 14, 17, 18, 19, 25</td>
</tr>
<tr>
<td>3.21</td>
<td>$1^+; 1$</td>
<td>$\approx 1000$</td>
<td>$n$</td>
<td>5, 14</td>
</tr>
<tr>
<td>5.4</td>
<td>$(2^+, 3^+); 1$</td>
<td>$\approx 650$</td>
<td>$n$</td>
<td>4, 5</td>
</tr>
<tr>
<td>6.1 ± 100</td>
<td></td>
<td>$\approx 900$</td>
<td>$n$</td>
<td>4</td>
</tr>
<tr>
<td>6.530 ± 20</td>
<td></td>
<td>$&lt; 40$</td>
<td>(n)</td>
<td>4, 19</td>
</tr>
<tr>
<td>7.1 ± 100</td>
<td></td>
<td>$\approx 350$</td>
<td>$n$</td>
<td>4</td>
</tr>
<tr>
<td>(9.)</td>
<td></td>
<td>$\approx 6000$</td>
<td></td>
<td>14, 17</td>
</tr>
</tbody>
</table>


$J = 2$ (1973NE10);

$\mu = +1.65335 \pm 0.00035$ nm (1973HA12);

$\mu = +1.6532 \pm 0.0008$ nm (1959CO68, 1962CO08);

$\mu = +1.6530 \pm 0.0008$ nm (1967GU14). See also (1967SH14, 1969FU11, 1971SH26).

1. $^8$Li($\beta^-$)$^8$Be $Q_m = 16.006$

The $\beta$-decay leads mainly to $^8$Be*(2.9): see $^8$Be. Measurements of the half-life of $^8$Li include $\tau_{1/2} = 841 \pm 4$ msec (1954KL36), 848 ± 5 msec (1960JA12), 844.0 ± 0.7 msec (1966CL02),
854 ± 8 msec \((1968DA12)\), 838 ± 6 msec \((1971WI05)\): \(\tau_{1/2} = 842 ± 6\) msec is adopted. [For other measurements see Table 8.2 in \((1966LA04)\).] See also \((1968BO32)\). Taking \(\tau_{1/2} = 842\) msec and \(Q = 16.006 - 2.94\), \(\log ft = 5.61\) (B. Zimmerman, private communication). See also \((1966BA1A, 1969BA43)\). The distribution of recoil momenta indicates \(J^p = 2^+\) (see \(^8\)Be).

The coefficient for the angular correlation between \(\beta^-\), \(\bar{\nu}_e\) and \(\alpha\)-particles, \(b = -0.88 ± 0.08\) \((1965GR25)\), \(-1.01 ± 0.07\) \((1966EII02)\), in substantial agreement with \(b = -1\), expected from axial vector coupling.

See also \((1970SC34)\) and \(^8\)Be.

2. \(^6\)Li\((t, p)^8\)Li

\(Q_m = 0.801\)

Transitions to \(^8\)Li\(^*\)(0, 0.98) have been observed: see \((1966LA04)\). See also \((1968CO1H)\).

3. \(^7\)Li\((n, \gamma)^8\)Li

\(Q_m = 2.0327\)

\(Q_0 = 2032.78 ± 0.15\) keV (E.T. Jurney, private communication);

\(Q_0 = 2032.8 ± 1\) keV \((1967RA24)\).

The thermal capture cross section is \(45.4 ± 3\) mb. Neutron capture \(\gamma\)-rays are observed with \(E_\gamma = 980.6 ± 0.2, 1052.0 ± 0.2\) and \(2032.5 ± 0.28\) keV, with intensities of \(10.6 ± 1, 10.6 ± 1\) and \(89.4 ± 1 \gamma/100\) neutrons: \(E_x = 980.7 ± 0.2\) keV for the first excited state (E.T. Jurney, private communication). The cross section for capture radiation has been measured for \(E_n = 40\) to \(1000\) keV: it decreases from \(50\) \(\mu\)b to \(5\) \(\mu\)b over that interval. The cross section shows the resonance corresponding to \(^8\)Li\(^*\)(2.26): \(\Gamma_\gamma = 0.07 ± 0.03\) eV \((1959IM04)\). See also \((1967GU14), (1966LA04)\) and \((1970AU1B, 1973MU14)\); astrophys. calculations).

4. \(^7\)Li\((n, n)^7\)Li

\(E_b = 2.0327\)

The thermal cross section is \(1.07 ± 0.04\) b \((1960HU1A)\); the coherent scattering length (thermal, bound) is \(-2.1 ± 0.1\) fm \((1969BA1P, 1973MU14)\). Total cross-section measurements have recently been reported at \(E_n = 0.010\) to \(1.236\) MeV \((1968HI1E)\), \(0.10\) to \(1.50\) MeV \((1970ME1C)\), \(1.12\) to \(2.30\) MeV \((1968KN1B)\), \(2.5\) to \(15.0\) MeV \((1971FO1A)\), \(3.35, 4.83, 5.74\) and \(7.5\) MeV \((1968HO03; also \sigma_{el})\), \(10\) MeV \((1967CO01)\) and \(0.6\) to \(30.0\) MeV (C.A. Goulding, private communication, and \((1973GO2B)\)). See also \((1968EN1A, 1972PR03)\). Polarization measurements at \(E_n = 4.4\) MeV are reported by \((1966ST09; n_0)\). For earlier references see \((1966LA04)\). See also \((1966DA1B, 1972LA1F)\) and \((1966AG1A, 1966SE1E, 1967BE1F, 1967HO1F; theor.)\). For angular distribution measurements see \(^7\)Li and \((1970GA1A)\).
Table 8.2: Resonance parameters for $^8$Li*(2.26) (1970ME1C) a,b

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E_{\text{res}}$ (keV)</td>
<td>261.2</td>
</tr>
<tr>
<td>$E_x$ (MeV) c</td>
<td>2.261</td>
</tr>
<tr>
<td>$\Gamma$ (keV)</td>
<td>35 ± 5 d</td>
</tr>
<tr>
<td>$\Gamma_n(E_r)$ (keV)</td>
<td>36.5</td>
</tr>
<tr>
<td>$\Gamma_\gamma$ (eV) c</td>
<td>0.07 ± 0.03 c</td>
</tr>
<tr>
<td>$\gamma_0^2$ (keV)</td>
<td>594</td>
</tr>
<tr>
<td>$\theta^2$</td>
<td>0.091</td>
</tr>
<tr>
<td>radius (fm)</td>
<td>3.30</td>
</tr>
<tr>
<td>$\sigma_{\text{max}}$ (b)</td>
<td>12.0 f</td>
</tr>
<tr>
<td>$J^\pi$</td>
<td>3$^+$</td>
</tr>
<tr>
<td>$l_n$</td>
<td>1</td>
</tr>
</tbody>
</table>

a See also (1956TH06, 1956WI04, 1960HU1A).

b Energies in laboratory system except for those labeled c.

c Energies in c.m. system.

d (1960HU1A).

e (1959IM04).

f See also (1968HI1E).

A pronounced resonance is observed at $E_n = 261$ keV with $J^\pi = 3^+$, formed by p-waves (Table 8.2) (1970ME1C). A good account of the polarization is given by the assumption of levels at $E_n = 0.25$ and 3.4 MeV, with $J^\pi = 3^+$ and $2^-$, together with a broad $J^\pi = 3^-$ level at higher energy (1964LA19). Broad peaks are reported at $E_n = 4.6$ and 5.8 MeV ($\pm 0.1$ MeV) [$^8$Li*(6.1, 7.1)] with $\Gamma \approx 1.0$ and 0.4 MeV, respectively, and there is indication of a narrow peak at $E_n = 5.1$ MeV [$^8$Li*(6.5)] with $\Gamma \ll 80$ keV and of a weak, broad peak at $E_n = 3.7$ MeV (1971FO1A: D.G. Foster and C.A. Goulding, private communications). See also reaction 5.

5. $^7$Li(n, n')$^7$Li* $E_b = 2.0327$

The excitation function for 0.48 MeV $\gamma$-rays shows an abrupt rise from threshold (indicating s-wave formation and emission) and a broad maximum ($\Gamma \approx 1$ MeV) at $E_n = 1.35$ MeV. A good fit is obtained with either $J^\pi = 1^-$ or 1$^+$ (2$^+$ not excluded), $\Gamma_{\text{lab}} = 1.14$ MeV (1955FR10). At higher energies a prominent peak is observed at $E_n = 3.8$ MeV ($\Gamma_{\text{lab}} = 0.75$ MeV) and there is
some indication of a broad resonance ($\Gamma_{\text{lab}} = 1.30 \text{ MeV}$) at $E_n = 5.0 \text{ MeV}$. Between $E_n = 6.0$ and 9.0 MeV, the cross section decreases monotonically. The two resonances are interpreted as being due to states with $J^{\pi} = 2^+$ or $3^+$ [$^9\text{Li}^*(5.4)$] and $J^{\pi} = 3^+$ or $3^-$ [$^9\text{Li}^*(6.4)$] [(1972PR03: $\Gamma_{\lambda}$, $\gamma_{\lambda}$ and $\theta_{\lambda}$ are listed for these two states under various assumptions)]. Cross-section measurements have also recently been reported for $E_n = 3.35, 4.83, 5.74$ and 7.5 MeV (1968HO03; $n_1$, $n_2$), 10 MeV (1967CO01: $n_2$) and 19.0 to 21.0 MeV (1972PR03; $\gamma_1$). See also (1965DE1G, 1972CO1K), and (1966LA04) for a listing of earlier references.

6. $^7\text{Li}(n, 2n)^6\text{Li}$ 
   $Q_m = -7.251$ 
   $E_b = 2.0327$


7. (a) $^7\text{Li}(n, p)^7\text{He}$ 
   $Q_m = -10.42$ 
   $E_b = 2.0327$

   (b) $^7\text{Li}(n, p)^4\text{He} + 3n$

   At $E_n = 19 \text{ MeV}$, the upper limit for reaction (b) is 10 mb (1971KO24). See also $^7\text{He}$ and (1973LI02).

8. $^7\text{Li}(n, d)^6\text{He}$ 
   $Q_m = -7.753$ 
   $E_b = 2.0327$

   See (1967VA12, 1973LI02) and (1966LA04).

9. $^7\text{Li}(n, t)^4\text{He} + n$ 
   $Q_m = 2.4668$ 
   $E_b = 2.0327$

   The cross section rises to 450 mb at $E_n \approx 8 \text{ MeV}$ and thereafter decreases slowly to 300 mb at $E_n = 15 \text{ MeV}$ (1964ST25). The large cross section, comparable to the geometric value, is understood in terms of the ($\alpha + t$) cluster nature of $^7\text{Li}$ (1962RO12). Cross sections for this reaction have recently been reported at $E_n = 3.35, 4.83, 5.74$ and 7.5 MeV (1968HO03) and at 10 MeV (1967CO01). See also (1967VA12, 1971AN1M, 1973LI02) and $^3\text{He}$, and (1966JE1B).

10. $^7\text{Li}(d, p)^8\text{Li}$ 
    $Q_m = -0.1919$ 
    $Q_0 = -188 \pm 7 \text{ keV}$ (1967SP09).
Proton groups have been observed to \(^8\)Li\(^*(0, 0.98, 2.26)\): see (1966LA04). At \(E_d = 15\) MeV, no states are observed in the region \(2.3 < E_x < 8\) MeV: a limit of 0.6 mb/sr is placed on groups with widths \(\lesssim 100\) keV \((\theta = 10^\circ, 14^\circ, 25^\circ)\) (1960HA14). Angular distributions of the \(p_0\) and \(p_1\) groups \([l_n = 1]\) at \(E_d = 12\) MeV have been analyzed by DWBA: \(S_{\text{exp.}} = 0.87\) and 0.48, respectively for \(^8\)Li\(^*(0, 0.98)\) \([S_{\text{exp.}}/S_{\text{theor.}} = 0.84, 1.09]\) (1967SC29). For PWBA analyses, see (1966LA04).

The first excited state decays by emitting a \(980 \pm 10\) keV \(\gamma\)-ray. The transition is M1 (1962CH14). The lifetime of \(^8\)Li\(^*(0.98)\) is \(10.1 \pm 4.5\) fsec (1971TH02). See also (1966TH1B, 1968FI1F).

11. \(^7\)Li(t, d)\(^8\)Li \[Q_m = -4.2249\]

See (1970CH1Q) and \(^{10}\)Be. See also (1970JA1J; theor.).

12. \(^7\)Li(\(\alpha\), \(^3\)He)\(^8\)Li \[Q_m = -18.546\]

Not reported.

13. \(^8\)He(\(\beta^-\))\(^8\)Li \[Q_m = 10.70\]

\(^8\)He decays with \(\tau_{1/2} = 122 \pm 2\) msec to \(^8\)Li\(^*(0.98)\) with an 88\% branch \([\log ft = 4.20; B. Zimmerman, private communication]\) and to neutron unstable states with a 12 \(\pm\) 1\% branch. The allowed decay supports the assignment \(J^\pi = 1^+\) to \(^8\)Li\(^*(0.98)\) (1965PO06).

14. \(^9\)Be(\(\gamma\), p)\(^8\)Li \[Q_m = -16.888\]


15. \(^9\)Be(e, ep)\(^8\)Li \[Q_m = -16.888\]

See (1969BA1F; theor.). See also \(^9\)Be.

16. \(^9\)Be(n, d)\(^8\)Li \[Q_m = -14.664\]
See (1969SC05) in $^{10}$Be.

17. $^{9}$Be(p, 2p)$^8$Li \[ Q_m = -16.888 \]

The summed proton spectrum at $E_p = 156$ MeV shows peaks corresponding to $^8$Li(0) and $^8$Li*(0.98 + 2.26) [unresolved]. In addition s-states [$J^\pi = 1^-, 2^-$] are suggested at $E_x = 9$ and 16 MeV, with $\Gamma_{c.m.} \approx 6$ and 8 MeV: the latter may actually be due to continuum protons (1967RO06). See also $^9$Be and (1966LA04) [(1965TY1A) has now been published as (1966TY01)]. See also (1971GR1K), (1965CO1E) and (1967KO1B, 1968JA1D, 1973AS02; theor.).

18. $^{9}$Be(d, $^3$He)$^8$Li \[ Q_m = -11.395 \]

At $E_d = 38$ MeV, differential cross sections have been obtained for the $^3$He groups corresponding to $^8$Li*(0, 0.98, 2.26) (1966GA21: see $^9$Be(d, t)$^8$Be).

19. $^{9}$Be(t, $\alpha$)$^8$Li \[ Q_m = 2.926 \]

At $E_t = 12.98$ MeV, $\alpha$-particle groups are observed to $^8$Li*(0, 0.98, 2.26) and to a state at $E_x = 6.530 \pm 0.020$ MeV with $\Gamma_{c.m.} < 40$ keV (1965WA12). For angular distributions to these four states, see (1968AJ01): the distribution of the $\alpha$-particles to $^8$Li*(6.53) is rather featureless and does not involve a forward maximum, suggesting $l > 1$. A large $l$-transfer is consistent with the narrow width of this unbound state (1968AJ01). At $E_t = 20$ MeV these four $^8$Li states were also observed. $^8$Li*(6.53) was found to have $\Gamma_{c.m.} = 40 \pm 10$ keV. No other groups corresponding to sharp states of $^8$Li ($\Gamma \lesssim 100$ keV) with $E_x \lesssim 16$ MeV were observed at $\theta = 15^\circ$, 25$^\circ$ and 35$^\circ$ (F. Ajzenberg-Selove and O. Hansen, private communication). The $\alpha_0$ angular distribution has also been measured at $E_t = 2.10$ MeV (1970CO04) and 0.52 to 1.31 MeV (1969NA04; also $\alpha_1$ at $E_t = 0.52$ to 1.67 MeV). The mean lifetime, $\tau_m = 14 \pm 5$ fsec for $^8$Li*(0.98): $E_x = 980.80 \pm 0.10$ keV (1972CO09).

See also $^{12}$B in (1975AJ02).

20. $^{9}$Be($^7$Li, $^8$Be)$^8$Li \[ Q_m = 0.368 \]

See (1966LE10) and (1968TO1C; theor.).

21. $^{10}$Be(d, $\alpha$)$^8$Li \[ Q_m = 2.372 \]
This reaction has been observed for $E_\alpha = 0.7$ to 3.0 MeV: see $^{12}$B (1970GO11, 1973GO09).

22. $^{10}$B(n, $^3$He)$^8$Li

\[ Q_m = -15.755 \]

Not reported.

23. $^{11}$B(\gamma, $^3$He)$^8$Li

\[ Q_m = -27.211 \]

See (1963NE07).

24. $^{11}$B(n, $\alpha$)$^8$Li

\[ Q_m = -6.633 \]

Angular distributions have been obtained at $E_n = 14.1$ MeV for the $\alpha_0$ and $\alpha_1$ groups (1973BO26). See also (1966JA1C) and $^{12}$B in (1975AJ02).

25. $^{12}$C($\alpha$, $^8$B)$^8$Li

\[ Q_m = -41.445 \]

At $E_\alpha = 129$ MeV, $^8$B groups are observed to $^8$Li*(0, 0.98, 2.26) (1968MC02).

26. $^{13}$C($^7$Li, $^{12}$C)$^8$Li

\[ Q_m = -2.914 \]

At $E(\tilde{7}$Li) = 34 MeV angular distributions have been measured for the transitions to $^8$Li\textsubscript{g.s.} + $^{12}$C\textsubscript{g.s.}, $^8$Li\textsubscript{0.98} + $^{12}$C\textsubscript{g.s.}, $^8$Li\textsubscript{g.s.} + $^{12}$C\textsubscript{4.4} and $^8$Li\textsubscript{0.98} + $^{12}$C\textsubscript{4.4} (1973SC26).

27. $^{18}$O($^7$Li, $^{17}$O)$^8$Li

\[ Q_m = -6.014 \]

See (1969NE1E).
GENERAL: (See also (1966LA04).)


Table 8.3: Energy levels of $^8$Be

<table>
<thead>
<tr>
<th>$E_x$ (MeV $\pm$ keV)</th>
<th>$J^\pi; T$</th>
<th>$\Gamma_{c.m.}$</th>
<th>Deacy</th>
<th>Reactions</th>
</tr>
</thead>
<tbody>
<tr>
<td>g.s.</td>
<td>$0^+; 0$</td>
<td>$6.8 \pm 1.7$ eV</td>
<td>$\alpha$</td>
<td>1, 4, 12, 13, 14, 15, 22, 23, 24, 25, 27, 30, 31, 32, 33, 34, 36, 37, 38, 39, 40, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 60, 62, 64</td>
</tr>
<tr>
<td>$2.94 \pm 30^a$</td>
<td>$2^+; 0$</td>
<td>$1.56 \pm 0.03$ MeV</td>
<td>$\alpha$</td>
<td>4, 13, 14, 15, 22, 23, 24, 25, 27, 28, 29, 30, 31, 32, 33, 36, 37, 38, 39, 40, 42, 44, 45, 46, 47, 48, 49, 50, 51, 52, 55</td>
</tr>
<tr>
<td>$(6.0 \pm 3000)$</td>
<td>$0^+; 0$</td>
<td>$(9 \pm 4$ MeV)</td>
<td>$\alpha$</td>
<td>4</td>
</tr>
<tr>
<td>$(10.0 \pm 3000)$</td>
<td>$2^+; 0$</td>
<td>$(12^{+4}_{-2}$ MeV)</td>
<td>$\alpha$</td>
<td>4, 29</td>
</tr>
<tr>
<td>$11.4 \pm 300$</td>
<td>$4^+; 0$</td>
<td>$\approx 3.5$ MeV $^d$</td>
<td>$\alpha$</td>
<td>4, 14, 22, 24, 31, 32, 33, 38, 42, 49, 51</td>
</tr>
<tr>
<td>$16.627 \pm 4^b$</td>
<td>$2^+; 0 + 1$</td>
<td>$107 \pm 3$ keV</td>
<td>$(\gamma), \alpha$</td>
<td>2, 4, 13, 15, 22, 23, 24, 28, 29, 33, 37, 38, 42, 46, 49</td>
</tr>
<tr>
<td>$16.911 \pm 4^b$</td>
<td>$2^+; 0 + 1$</td>
<td>$77 \pm 3$ keV</td>
<td>$(\gamma), \alpha$</td>
<td>2, 4, 13, 15, 22, 23, 24, 31, 32, 33, 37, 38, 42, 46, 49</td>
</tr>
<tr>
<td>$17.642 \pm 1.5^b,c$</td>
<td>$1^+; 1$</td>
<td>$10.7 \pm 0.5$ keV</td>
<td>$\gamma, p$</td>
<td>13, 16, 18, 22, 23, 31, 32, 33, 38, 42</td>
</tr>
<tr>
<td>$18.154 \pm 4^b,c$</td>
<td>$1^+; 0$</td>
<td>$138 \pm 6$ keV</td>
<td>$\gamma, p$</td>
<td>13, 16, 18, 22, 31, 33, 38, 42</td>
</tr>
<tr>
<td>$18.91^c$</td>
<td>$2^-; 0$</td>
<td>$48 \pm 20$ keV</td>
<td>$\gamma, n, p$</td>
<td>13, 16, 17, 18, 22, 31, 32, 42</td>
</tr>
<tr>
<td>$19.06 \pm 20^c$</td>
<td>$3^+$</td>
<td>$270 \pm 20$ keV</td>
<td>$\gamma, p$</td>
<td>13, 16, 18, 22, 31, 32, 42</td>
</tr>
<tr>
<td>$19.22$</td>
<td>$3^+; (1)$</td>
<td>$208 \pm 30$ keV</td>
<td>$n, p$</td>
<td>13, 17, 18, 22, 31, 33, 38</td>
</tr>
<tr>
<td>$19.4$</td>
<td>$1^-$</td>
<td>$\approx 650$ keV</td>
<td>$n, p$</td>
<td>13, 17, 18</td>
</tr>
<tr>
<td>$19.9$</td>
<td>$4^+; 0$</td>
<td>$&lt; 1$ MeV</td>
<td>$\alpha$</td>
<td>4, 13, 24, 32, 33</td>
</tr>
</tbody>
</table>
Table 8.3: Energy levels of $^8$Be (continued)

<table>
<thead>
<tr>
<th>$E_x$ (MeV ± keV)</th>
<th>$J^\pi; T$</th>
<th>$\Gamma_{c.m.}$ (keV)</th>
<th>Deacy</th>
<th>Reactions</th>
</tr>
</thead>
<tbody>
<tr>
<td>20.1</td>
<td>$2^+; 0$</td>
<td>$\approx 1.1$ MeV</td>
<td>n, p, $\alpha$</td>
<td>4, 17, 18, 21, 32</td>
</tr>
<tr>
<td>20.2</td>
<td>$0^+; 0$</td>
<td>$&lt; 1$ MeV</td>
<td>$\alpha$</td>
<td>4</td>
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<tr>
<td>20.9 ± 200</td>
<td>$4^-$</td>
<td>$1.6 \pm 0.2$ MeV</td>
<td>p</td>
<td>18</td>
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<tr>
<td>21.5 ± 300</td>
<td>(3$^+$)</td>
<td>$1$ MeV</td>
<td>$\gamma$, p, n</td>
<td>16, 17, 31</td>
</tr>
<tr>
<td>22.0</td>
<td>$1^-; 1$</td>
<td>$4 - 5$ MeV</td>
<td>$\gamma$, p</td>
<td>16</td>
</tr>
<tr>
<td>22.2</td>
<td>$2^+; 0$</td>
<td>$\approx 0.8$ MeV</td>
<td>n, p, d, $\alpha$</td>
<td>4, 6, 7, 11, 15, 18, 21, 33</td>
</tr>
<tr>
<td>23.6</td>
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<td>$\gamma$, p</td>
<td>16</td>
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<tr>
<td>24.0</td>
<td>$1^-, 2^-$</td>
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<td>$\gamma$, p</td>
<td>16</td>
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<tr>
<td>25.2</td>
<td>$2^+; 0$</td>
<td>$\approx 1$ MeV</td>
<td>p, d, $\alpha$</td>
<td>4, 7, 11, 21</td>
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<tr>
<td>25.5</td>
<td>$4^+; 0$</td>
<td></td>
<td>$\alpha$</td>
<td>4</td>
</tr>
<tr>
<td>27.483 ± 10</td>
<td>$0^+; 2$</td>
<td>$10 \pm 3$ keV</td>
<td>p, d, $\alpha$</td>
<td>7, 11</td>
</tr>
<tr>
<td>(28.6)</td>
<td></td>
<td>broad</td>
<td>$\gamma$, p</td>
<td>16</td>
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</tbody>
</table>

$^a$ See also Table 8.4.
$^b$ See also Table 8.5.
$^c$ See also Table 8.7.
$^d$ We are greatly indebted to Prof. F.C. Barker for enlightening discussions concerning the width of $^8$Be*(11.4).


Adjusted mass excess $^\dagger$ of $^8$Be: 4941.87 ± 0.13 keV (1972WA1G).

1. $^8$Be $\rightarrow$ $^4$He$^4$He $\quad Q_m = 0.09189$
$\quad Q_0 = 92.6 \pm 0.8$ keV (1966RE02).

In $\alpha$-$\alpha$ scattering (reaction 4) the $Q_0$ is found to be 92.12 ± 0.05 keV, $\Gamma_{c.m.} = 6.8 \pm 1.7$ eV (1968BE02). See also (1966LA04) for earlier values.

$^\dagger$ Not used in $Q_m$ calculations in this paper.
2. $^4\text{He}(\alpha, \gamma)^8\text{Be}$ 

$$Q_m = -0.09189$$

Radiative widths have been measured for $^8\text{Be}^*(16.6, 16.9)$: see (1974NA1H: $E_\alpha = 31$ to 35 MeV).

3. (a) $^4\text{He}(\alpha, n)^7\text{Be}$ 
(b) $^4\text{He}(\alpha, p)^7\text{Li}$

$$Q_m = -18.9921 \quad E_b = -0.09189$$

$$Q_m = -17.348$$

For reaction (a) see (1952WA31). For reaction (b) see $^7\text{Li}$. 

4. $^4\text{He}(\alpha, \alpha)^4\text{He}$ 

$$E_b = -0.09189$$

Alpha-$\alpha$ scattering reveals the ground state as a resonance with $Q_\theta = 92.12 \pm 0.05$ keV, $\Gamma_{c.m.} = 6.8 \pm 1.7$ eV, $[\tau = (0.97 \pm 0.24) \times 10^{-16}$ sec] (1966BE05, 1968BE02). Effective range theory analysis of higher energy scattering yields widths consistent with this value but subject to considerable uncertainty (1966TO1B). However, (1967KE1B, 1967KE1E) and (1967RA1B) have carried out such analyses yielding $\Gamma = 6.14 \pm 0.04$ eV and $6.4^{+0.8}_{-0.5}$ eV, respectively. Using a three-level, one-channel $R$-matrix formalism (1968BA2D) find $\Gamma = 5.1 \pm 0.4$ eV for the ground state of $^8\text{Be}$. $R$-matrix analysis of the s-wave scattering and of the $^8\text{Be}(p, d)^8\text{Be}$ reaction indicates the presence of a second 0$^+$ state at $E_x \approx 6 \pm 3$ MeV, $\Gamma = 9 \pm 4$ MeV ($a_0 = 7$ fm) (1968BA2D). For $E_\alpha = 30$ to 70 MeV the $l = 0$ phase shift shows resonant behavior at $E_\alpha = 40.7$ MeV, corresponding to a 0$^+$ state at $E_x = 20.2$ MeV, $\Gamma < 1$ MeV, $\Gamma_0/\Gamma < 0.5$. No evidence for other 0$^+$ states is seen above $E_\alpha = 43$ MeV (1972BA83).

The d-wave phase shift becomes appreciable for $E_\alpha > 2.5$ MeV and passes through resonance at $E_\alpha = 6$ MeV ($E_x = 3.18$ MeV, $\Gamma = 1.5$ MeV, $J^\pi = 2^+$) (1963TO02). See Table 8.4. Analyses by many-level $R$-matrix theory of the $\alpha$-scattering, of the $^8\text{Be}(p, d)^8\text{Be}$ reaction and of the $^8\text{Li}$ and $^8\text{B}$ $\beta$-decays lead to approximately correct values for the $E_x$ and $\Gamma$ of $^8\text{Be}^*(2.9)$ and suggest a second 2$^+$ state at $E_x \approx 8.5$ MeV, $\Gamma \approx 10.5$ MeV (1969BA43: $a_2 = 6.75$ fm), $E_x = 12.0^{+3.0}_{-3.5}$ MeV, $\Gamma = 14.5^{+7.5}_{-6}$ MeV (1969CL10: $a_2 = 6.0 \pm 0.5$ fm). Five 2$^+$ levels are observed from the $l = 2$ phase shifts measured from $E_\alpha = 30$ to 70 MeV. $^8\text{Be}^*(16.6, 16.9)$ with $\Gamma_0 = \Gamma$, and states with $E_x = 20.2, 22.2$ and 25.2 MeV. The latter has a small $\Gamma_0$ (1972BA83). See also reaction 11.

The $l = 4$ shift rises from $E_\alpha \approx 11$ MeV and indicates a broad 4$^+$ level at $E_x = 11.4 \pm 0.3$ MeV (1959BR71). See also (1973CH1W). A rapid rise of $\delta_4$ at $E_\alpha = 40$ MeV corresponds to a 4$^+$ state at 19.9 MeV with $\Gamma_0/\Gamma \approx 0.96$; $\Gamma < 1$ MeV and therefore $\Gamma_0 < 1$ MeV, which is < 5% of the Wigner limit. A broad 4$^+$ state is also observed near $E_\alpha = 51.3$ MeV ($E_x = 25.5$ MeV) but there is no evidence for a previously reported state at $E_x = 27.5$ MeV (1972BA83).

Over the range $E_\alpha = 30$ to 70 MeV a gradual increase in $\delta_6$ is observed (1972BA83). Some indications of a 6$^+$ state at $E_x \approx 28$ MeV and of an 8$^+$ state at $\approx 57$ MeV have been reported.
Table 8.4: Parameters of $^8\text{Be}^*(2.9)$

<table>
<thead>
<tr>
<th>$E_x$ (MeV)</th>
<th>$\Gamma_{\text{cm}}$ (MeV)</th>
<th>Reaction</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>2.83 ± 0.20</td>
<td>1.75 ± 0.30</td>
<td>$^6\text{Li}(^3\text{He}, \text{p}), ^{10}\text{B}(\alpha, \text{p})$</td>
<td>(1969NU01)</td>
</tr>
<tr>
<td>3.1 ± 0.1</td>
<td>1.2 ± 0.3</td>
<td>$^6\text{Li}(\alpha, \text{d})$</td>
<td>(1969BA18)</td>
</tr>
<tr>
<td>3.10 ± 0.09</td>
<td>1.75 ± 0.1</td>
<td>$^7\text{Li}(\text{d}, \text{n})$</td>
<td>(1964JO04)</td>
</tr>
<tr>
<td>2.90 ± 0.06</td>
<td>1.53 ± 0.04</td>
<td>$^7\text{Be}(\text{d}, \text{p})$</td>
<td>(1960KA17)</td>
</tr>
<tr>
<td>2.90 ± 0.04</td>
<td>1.35 ± 0.15</td>
<td>$^9\text{Be}(^3\text{He}, \alpha)$</td>
<td>(1963DO08)</td>
</tr>
<tr>
<td></td>
<td>1.48 ± 0.07</td>
<td>$^{11}\text{B}(\text{p}, \alpha)$</td>
<td>(1971KA21)</td>
</tr>
<tr>
<td>2.94 ± 0.03</td>
<td>1.56 ± 0.03</td>
<td>mean</td>
<td></td>
</tr>
</tbody>
</table>

$^a$ See also Table 8.9 in (1966LA04) and reaction 32 (1973SO08).

by (1965DA1A), with $\Gamma_{\text{cm}} \approx 20$ and $\approx 73$ MeV, respectively. The elastic scattering has been studied at $E_\alpha = 140$ MeV by (1972FR1K). For a listing of the older work see Table 8.7 [(d$\sigma$/d$\Omega$)], Table 8.8 [parameters of $^8\text{Be}$ states from $^4\text{He}(\alpha, \alpha)$] and Table 8.9 [parameters of $^8\text{Be}^*(2.9)$] in (1966LA04). See also (1972FR1J). For studies of inelastic scattering of $\alpha$-particles from $^4\text{He}$ see (1969GR06, 1971HA41) and the review in (1973FI04). See also (1965SL1A, 1967ST30, 1968CO1M).

The bremsstrahlung cross section has been measured at $E_\alpha = 9.35$ MeV and for $E_\alpha = 11.4$ to 13.5 MeV: no significant enhancement is found at the final state energy corresponding to $^8\text{Be}^*(2.9)$ (1972FR02, 1973FR17). The cross section has also been measured for $E_\alpha = 12.1$ to 18.7 MeV by (1972PE16).


5. $^6\text{Li}(\gamma)^8\text{Be}$

\[ Q_m = 22.282 \]

Not observed: (1953SA1A, 1954SI07).
6. (a) $^6\text{Li}(d, n)^7\text{Be}$

$$Q_m = 3.382$$

$$E_b = 22.282$$

(b) $^6\text{Li}(d, n)^4\text{He} + ^3\text{He}$

$$Q_m = 1.795$$

The yield curve has been measured for $E_d = 0.06$ to 5.5 MeV [see (1966LA04) and (1966SC26)], and at $E_d = 12$ to 17 MeV (1970GA07: population of $^7\text{Be}^*$(0 + 0.43)). A broad s-wave resonance is indicated at $E_d = 0.41$ MeV, $\Gamma = 0.45$ MeV (1952BA1A, 1956NE13). Polarization measurements are reported at $E_d = 0.27$ to 0.60 MeV (1966MI06; n$_1$) and 2.5 to 3.7 MeV (1970TH08; n$_0$, n$_1$). The distributions observed by (1970TH08) are quite constant over the range 2.5 to 3.7 MeV, indicating that the predominant reaction mechanism is direct. See also $^7\text{Be}$.

Comparisons of the populations of $^7\text{Be}^*$(0, 0.43) and of $^7\text{Li}^*$(0, 0.48) have been made at many energies, up to $E_d = 7.2$ MeV. The n/p ratios are closely equal for analog states, consistent with charge symmetry (1957WI24, 1963BI1B, 1963CR08). See also (1966AU1A).

7. (a) $^6\text{Li}(d, p)^7\text{Li}$

$$Q_m = 5.026$$

$$E_b = 22.282$$

(b) $^6\text{Li}(d, p)^4\text{He} + ^3\text{H}$

$$Q_m = 2.559$$

Excitation functions have been measured for $E_d = 30$ keV to 5.4 MeV [see (1966LA04) and (1966LO18, 1968BE1P, 1969BL14)]. A broad maximum near $E_d = 1.0$ MeV is interpreted as indicating a level at $E_d = 0.4$ MeV (1950WH02). In the range $E_d = 1$ to 5 MeV there is evidence for both direct interaction and compound nucleus formation (1963BI1B, 1963ME09, 1964PA06): at back angles the (d, p$_1$) data show evidence of the $E_d = 3.7$ MeV resonance [see $^6\text{Li}(d, \alpha)^4\text{He}$]. An anomaly is observed in the p$_1$/p$_0$ intensity ratio at $E_d = 6.945$ MeV, corresponding to the $J^\pi = 0^+; T = 2$ analog of the ground state of $^8\text{He}$: $E_x = 27.483 \pm 0.010$ MeV, $\Gamma = 10 \pm 3$ keV, $\Gamma_{p_0} \ll \Gamma_{p_1}$, $\Gamma_{p_0} < \Gamma_d$ (1969BL14). The parameters of this state have been calculated by (1969BA36).

Polarization measurements have been made at $E_d = 0.6$ and 0.96 MeV (1969NA1J, 1972SE09; p$_0$, p$_1$) and at $E_d = 2.1$ to 10.9 MeV (1968DU09, 1968FI07, 1970FI07; p$_0$, p$_1$). The latter report pronounced differences in the angular distributions of the vector analyzing power of the two $l_n = 1$ transitions to $^7\text{Li}^*(0, 0.48)$. See also (1972FI1E, 1973FI1C). For reaction (b) see (1966FR06). See also $^7\text{Li}$ and (1966AU1A, 1966BR25).

8. $^6\text{Li}(d, d)^6\text{Li}$

$$E_b = 22.282$$

The yield of elastically scattered deuterons has been measured for $E_d = 2$ to 4.8 MeV (1964PA06), 4.0 to 6.5 MeV (1966BR1J), and 6.33 to 7.14 MeV (1969BL14): no resonances are reported. At $E_d = 12.0$ MeV, $\theta_{\text{lab}} = 95^\circ$, the differential cross section for elastic scattering is $9.82 \pm 0.20$ mb/sr (1971BI11). See also $^6\text{Li}$ and (1972FI1E, 1973FI1C).
9. $^6\text{Li}(d, t)^{^6}\text{Li}$  

$Q_m = 0.59 

E_b = 22.282$

The cross section for tritium production rises rapidly to 190 mb at 1 MeV, then more slowly to 290 mb near 4 MeV. There is evidence of deviation from isotropy near 0.4 MeV (1955MA20). See also $^5\text{Li}$.

10. $^6\text{Li}(d, \, ^3\text{He})^{^5}\text{He}$  

$Q_m = 0.90 

E_b = 22.282$

See $^5\text{He}$.

11. (a) $^6\text{Li}(d, \alpha)^4\text{He}$  

$Q_m = 22.374 

E_b = 22.282$

(b) $^6\text{Li}(d, \alpha p)^3\text{H}$  

$Q_m = 2.559$

(c) $^6\text{Li}(d, \alpha n)^3\text{He}$  

$Q_m = 1.795$

(d) $^6\text{Li}(d, 2d)^4\text{He}$  

$Q_m = -1.4737$

Cross sections and angular distributions (reaction (a)) have been measured for $E_d = 0.03$ to 12 MeV: see (1966LA04) and (1969LE22: $E_d = 40$ to 130 keV), (1966LO18: 0.2 to 2 MeV), (1968BE1P: 0.3 to 1.0 MeV), (1965RO1E, 1969DE31, 1969HU17: 0.35 to 1.50 MeV), (1966BR25: 1.0 to 2.0 MeV), (1964AN1A: 1.0 to 2.5 MeV), (1967CL06: 3 to 12 MeV) and (1969BL14: 6.33 to 7.14 MeV). Polarization measurements are reported at $E_d = 0.40, 0.60, 0.80, 0.96$ MeV (1971NE12, 1972SE09), 0.7 to 2.2 MeV (1967PL02), 2.1 to 10.9 MeV (1968DU07), 4.3, 6.3, 8.0, 10.1 and 11.8 MeV (1968BU13) and at $E(6\text{Li}) = 0.6$ MeV (1970HO11). See also (1967BU1B, 1972KO1P).

Maxima are observed at $E_d = 0.8$ MeV, $\Gamma_{lab} \approx 0.8$ MeV and $E_d = 3.75$ MeV, $\Gamma_{lab} \approx 1.4$ MeV (1963ME09, 1964PA06). Analysis of these and other data up to $E_d = 12$ MeV indicate a $2^+$, $0^+$, $(6^+)$, $2^+$, $4^+$ sequence of states: see Table 8.10 in (1966LA04) (1965FR02, 1967CL06). See, however, reaction 4.

The assignment of $J^\pi = 2^+$ to $^8\text{Be}^*$ (22.2) is consistent with the polarization information (1971NE12, 1972SE09), but the $0^+$ state may actually be virtual with respect to $^6\text{Li}+d$ (1972SE09) [$^8\text{Be}^*$ (20.3); see reaction 4 and (1972BA83)]. At $E_d = 6.945$ MeV, the $\alpha$-yield shows an anomaly corresponding to $^8\text{Be}^*$ (27.48), the $J^\pi = 0^+$; $T = 2$ analog of the $^8\text{He}$ ground state (1969BL14). See also reaction 7 and (1969BA36; theor.).

See also (1966LE1C), (1971PL1C) and (1967TS1A, 1968CO1L, 1968KO1G, 1969CH1J, 1970FI11; theor.). For reactions (b), (c) and (d), see (1972HA34) and (1973FI04).

12. (a) $^6\text{Li}(t, n)^8\text{Be}$  

$Q_m = 16.024$

(b) $^6\text{Li}(t, n)^4\text{He}^4\text{He}$  

$Q_m = 16.116$
For reaction (b) see (1966LA02, 1967BE13, 1967BI1D). See also (1966LA04).

13. (a) $^6$Li($^3$He, p)$^8$Be
   \[ Q_m = 16.788 \]
   (b) $^6$Li($^3$He, p)$^4$He
   \[ Q_m = 16.880 \]

Proton groups are observed to $^8$Be*(0, 2.9, 16.63, 16.91, 17.64): see (1966LA04) and Tables 8.4 and 8.5. The excitation of $^8$Be*(18.15, 19.0, 19.4, 19.9) is also reported by (1971GI07). Angular distributions have been measured at $E(^3$He) = 1.4 to 1.8 MeV (1969VI05; $p_0$, $p_1$) and 5, 6, 7, 9, 10, 13 and 17 MeV (1965FL03; $p_0$, $p_1$; PWBAE analysis). A gradual change is observed from a dominant back angle maximum to a dominant forward maximum (1965FL03). Measurements of the energies of all the particles emitted in this reaction and reactions 23, 38 and 42 show that the apparent width of $^8$Be*(2.9) does not depend on the relative velocity of the spectator particle: $E_x = 2.83 \pm 0.20$ MeV, $\Gamma = 1.75 \pm 0.30$ MeV (1969NU01). See also Table 8.4.

Reaction (b) proceeds via $^8$Be*(16.63, 16.91): $\Gamma = 117 \pm 10$ and $85 \pm 10$ keV, respectively. Interference effects are observed (1969VI05). See also $^6$Li and (1967RE03, 1968RE10, 1972TH08). See also (1964MA57, 1968VI03, 1970GA1G, 1971TR1B), (1967HO1C) and (1967BA1E, 1968HE1F, 1969TH1D, 1970DE41, 1973ED02; theor.).

14. (a) $^6$Li($\alpha$, d)$^8$Be
   \[ Q_m = -1.5656 \]
   (b) $^6$Li($\alpha$, 2$\alpha$)$^2$H
   \[ Q_m = -1.4737 \]

Deuteron groups have been observed to $^8$Be*(0, 2.9, 11.3 ± 0.4) (1959ZE1A, 1962CE01). Angular distributions have been measured at $E_\alpha = 20$ and 24 MeV (1973GR1N), 20.5 to 24.5 MeV and at 38 MeV (1965DE1F; $d_0$), 43 MeV (1959ZE1A; $d_0$, $d_1$) and 48 MeV (1962CE01; $d_0$, $d_1$). At $E_\alpha = 12$ MeV ($\theta = 15^\circ$ and $20^\circ$) the deuteron spectrum does not show a “ghost” anomaly at $E_x = 0.1 - 0.5$ MeV (1971BE52). A study of reaction (b) shows that the peak due to $^8$Be*(2.9) is best fitted by using $\Gamma = 1.2 \pm 0.3$ MeV (1969BA18): see also Table 8.4. See also (1968LA1E) and (1971BU1K; theor.). For reaction (b) see $^6$Li. See also (1966LA04).

15. (a) $^6$Li($^6$Li, $\alpha$)$^8$Be
   \[ Q_m = 20.808 \]
   (b) $^6$Li($^6$Li, $\alpha$)$^4$He$^4$He
   \[ Q_m = 20.900 \]

This reaction proceeds via $^8$Be*(0, 2.9, 16.6, 16.9, 22.5), and there is indication also that the direct three-body break-up (reaction (b)) is possible (1971GA1N, 1971GA21, 1972GA32: $E_{\text{max}}(^6$Li) = 13.0 MeV). The involvement of a state at $E_x = 19.9$ MeV ($\Gamma = 1.3$ MeV) is suggested by (1966MA40). See also (1971GI07). Good agreement with the shapes of the peaks
<table>
<thead>
<tr>
<th>$E_x$ (MeV ± keV)</th>
<th>$\Gamma_{c.m.}$ (keV)</th>
<th>Reaction</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>16.625 ± 10</td>
<td>95 ± 20</td>
<td>$^6$Li($^3$He, p)</td>
<td>(1961ER01)</td>
</tr>
<tr>
<td></td>
<td>117 ± 10</td>
<td>$^6$Li($^3$He, p)</td>
<td>(1969VI05)</td>
</tr>
<tr>
<td>16.627 ± 5</td>
<td>113 ± 3</td>
<td>$^7$Li($^3$He, d)</td>
<td>(1967MA12)</td>
</tr>
<tr>
<td>16.627 ± 15</td>
<td>105 ± 30</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
<td>(1961ER01)</td>
</tr>
<tr>
<td>16.635 ± 15</td>
<td>96 ± 20</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
<td>(1963DO08)</td>
</tr>
<tr>
<td>16.623 ± 10</td>
<td>95 ± 20</td>
<td>$^{10}$B(d, $\alpha$)</td>
<td>(1961ER01)</td>
</tr>
<tr>
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<td>90 ± 5</td>
<td>$^{10}$B(d, $\alpha$)</td>
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<td>16.627 ± 4</td>
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<td>16.931 ± 10</td>
<td>85 ± 20</td>
<td>$^6$Li($^3$He, p)</td>
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<td>85 ± 10</td>
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<td>16.901 ± 5</td>
<td>77 ± 3</td>
<td>$^7$Li($^3$He, d)</td>
<td>(1967MA12)</td>
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<td>103 ± 15</td>
<td>$^9$Be(p, d)</td>
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<td>16.914 ± 12</td>
<td>88 ± 25</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
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</tr>
<tr>
<td>16.930 ± 15</td>
<td>80 ± 15</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
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</tr>
<tr>
<td>16.919 ± 10</td>
<td>85 ± 20</td>
<td>$^{10}$B(d, $\alpha$)</td>
<td>(1961ER01)</td>
</tr>
<tr>
<td></td>
<td>70 ± 5</td>
<td>$^{10}$B(d, $\alpha$)</td>
<td>(1971NO04)</td>
</tr>
<tr>
<td>16.911 ± 4</td>
<td>77 ± 3</td>
<td>mean</td>
<td></td>
</tr>
<tr>
<td>17.642 ± 10</td>
<td>&lt; 20</td>
<td>$^6$Li($^3$He, p)</td>
<td>(1961ER01)</td>
</tr>
<tr>
<td>17.642 ± 1.5</td>
<td>10.7 ± 0.5</td>
<td>$^7$Li(p, $\gamma$)</td>
<td>Table 8.6</td>
</tr>
<tr>
<td>17.636 ± 10</td>
<td>&lt; 15</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
<td>(1961ER01)</td>
</tr>
<tr>
<td>17.641 ± 10</td>
<td>10.7 ± 0.5</td>
<td>$^9$Be($^3$He, $\alpha$)</td>
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<tr>
<td></td>
<td>best</td>
<td>best</td>
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</tr>
<tr>
<td>18.157 ± 5</td>
<td>147</td>
<td>$^7$Li(p, $\gamma$)</td>
<td>Table 8.6</td>
</tr>
<tr>
<td>18.150 ± 5 c</td>
<td>138 ± 6</td>
<td>$^{10}$B(d, $\alpha$)</td>
<td>(1970CA12)</td>
</tr>
<tr>
<td>18.154 ± 4</td>
<td>138 ± 6</td>
<td>mean</td>
<td></td>
</tr>
</tbody>
</table>

a See also Table 8.11 in (1966LA04).
b Based on listed $Q_m$.
c Based on $E_x = 17.642$ MeV.
corresponding to $^8\text{Be}^*(16.6, 16.9)$ is obtained by using a simple two-level formula with interference, corrected for the effect of final state Coulomb interactions, assuming $\Gamma(16.6) = 90$ and $\Gamma(16.9) = 70$ keV; see also Table 8.5 (1971NO04). See also (1968NO03, 1969IN06). The ratio of the intensities of the groups corresponding to $^8\text{Be}^*(16.6, 16.9)$ remains constant for $E(^6\text{Li}) = 4.3$ to 5.5 MeV: $I(16.6)/I(16.9) = 1.22 \pm 0.08$ (1966KI09, 1966MA40). Partial angular distributions for the $\alpha_0$ group have been measured at fourteen energies for $E(^6\text{Li}) = 4$ to 24 MeV (1970FR06). For reaction (b) see also (1966BE22).


16. $^7\text{Li}(p, \gamma)^8\text{Be}$ $Q_m = 17.256$

Cross sections and angular distributions have been reported from $E_p = 30$ keV to 18 MeV. Gamma rays are observed to the ground ($\gamma_0$) and to the broad, $2^+$, excited state at 2.9 MeV ($\gamma_1$) and to $^8\text{Be}^*(16.6, 16.9)$ ($\gamma_3, \gamma_4$). Resonances for both $\gamma_0$ and $\gamma_1$ occur at $E_p = 0.44$ and 1.03 MeV, and for $\gamma_1$ alone at 2, 4.9, 6.0, 7.3, and possibly at 3.1 and 11.1 MeV. In addition broad resonances are reported at $E_p \approx 5$ MeV ($\gamma_0$), $\Gamma \approx 4 - 5$ MeV, and at $E_p \approx 7.3$ MeV ($\gamma_1$), $\Gamma \approx 8$ MeV. The $E_p \approx 5$ MeV resonance ($E_x \approx 22$ MeV) represents the giant dipole resonance based on $^8\text{Be}(0)$ while the $\gamma_1$ resonance, $\approx 2.3$ MeV higher, is based on $^8\text{Be}^*(2.9)$. The $\gamma_0$ and $\gamma_1$ giant resonance peaks each contain about 10% of the dipole sum strength (1966FI1B, 1968BL1E, 1970FI1B). The main trend between $E_p = 8$ and 17.5 MeV is a decreasing cross section (1970FI1B). See, however, (1967FE04). See also Table 8.6.

At the $E_p = 0.44$ MeV resonance ($E_x = 17.64$ MeV) the radiation is nearly isotropic consistent with p-wave formation, $J^\pi = 1^+$, with channel spin ratio $\sigma(J_c = 2)/\sigma(J_c = 1) = 3.2 \pm 0.5$ (1961ME10). Radiative widths for the $\gamma_0$ and $\gamma_1$ decay are displayed in Table 8.7. The E2/M1 amplitude ratio for the 17.6 $\rightarrow$ 2.9 transition varies over the energy of the broad final state: the average value is $\delta = 0.21 \pm 0.04$ (1967CO19). See also (1967CO29).

$^8\text{Be}^*(16.63, 16.91)$ are $2^+$ states with mixed $T = 0$, 1 isospin [see (1965MA1G, 1966MA03, 1968PA09, 1969SW01)], with the lower state of $^7\text{Li} + p$ parentage and the higher of $^7\text{Be} + n$ parentage (1965SW03, 1968PA09). A careful study of the $\alpha$-breakup of $^8\text{Be}^*(16.63, 16.91)$ for $E_p = 0.44$ to 2.45 MeV shows that the non-resonant part of the cross section for production of $^8\text{Be}^*(16.63)$ is accounted for by an extranuclear direct-capture process. Resonances for production of $^8\text{Be}^*(16.63, 16.91)$ are observed at $E_p = 0.44, 1.03$ and 1.89 MeV [$^8\text{Be}^*(17.64, 18.15, 18.9)$]. The results are consistent with the hypothesis of nearly maximal isospin mixing for $^8\text{Be}^*(16.63, 16.91)$; decay to these states is not observed from the $3^+$ states at $E_x = 19$ MeV, but rather from the $2^-$ state at 18.9 MeV excitation (1969SW01). See also reaction 17. (1968PA09) find squared $T = 1$ components of 40% and 60% in $^8\text{Be}^*(16.6, 16.9)$ and of 95% and 5% in $^8\text{Be}^*(17.6, 18.2)$. Gamma-$\alpha$ angular correlation measurements at $E_p = 0.44$ MeV show that the $17.64 \rightarrow 16.63 \gamma$ is nearly pure M1 $\delta(E2/M1) = -0.014 \pm 0.013$ (1969SW02). See also (1965SW03). Radiative widths are shown in Table 8.7.
Table 8.6: $^8$Be levels from $^7$Li(p, $\gamma$)$^8$Be

<table>
<thead>
<tr>
<th>$E_{\text{res}}$ (keV)</th>
<th>$\Gamma_{\text{lab}}$ (keV)</th>
<th>$^8$Be*</th>
<th>$l_p$</th>
<th>$J^\pi$</th>
<th>Res. $^c$ in</th>
<th>$\omega\Gamma_\gamma$ (eV)</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>441.4 ± 0.5 $^a$</td>
<td>12.2 ± 0.5</td>
<td>17.642</td>
<td>1</td>
<td>$1^+$</td>
<td>$\gamma_0, \gamma_1, \gamma_3, \gamma_4$</td>
<td>9.4</td>
<td>(1949FO18, 1956BU27, 1969SW01)</td>
</tr>
<tr>
<td>1030 ± 5</td>
<td>168</td>
<td>18.157</td>
<td>1</td>
<td>$1^+$</td>
<td>$\gamma_0, \gamma_1, \gamma_3, \gamma_4$</td>
<td>2</td>
<td>(1954KR06, 1960MA33, 1963RI09, 1970FI1B)</td>
</tr>
<tr>
<td>1890</td>
<td>150 ± 50</td>
<td>18.91</td>
<td></td>
<td>(2$^-$)</td>
<td>$\gamma_3, \gamma_4$</td>
<td>(1969SW01)</td>
<td></td>
</tr>
<tr>
<td>2060 ± 20</td>
<td>310 ± 20</td>
<td>19.06</td>
<td></td>
<td>$J = 1, 2, 3$</td>
<td>$\pi = (-)^b$</td>
<td>$\gamma_1$</td>
<td>(1957NE1A, 1963PE15, 1963RI09, 1967NI1A, 1970FI1B)</td>
</tr>
<tr>
<td>(3100)</td>
<td>(20.0)</td>
<td></td>
<td>(0)</td>
<td>$1^-$</td>
<td>$\gamma_0$</td>
<td>(1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>4900</td>
<td>21.5</td>
<td></td>
<td></td>
<td></td>
<td>$\gamma_1$</td>
<td>(1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>5000</td>
<td>≈ 4500</td>
<td>21.6</td>
<td>(0)</td>
<td>$1^-$</td>
<td>$\gamma_0$</td>
<td>(1959GE33, 1963MI08, 1963PE15, 1966FI1B, 1968BL1E, 1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>6000</td>
<td>22.5</td>
<td></td>
<td></td>
<td></td>
<td>$\gamma_1$</td>
<td>(1963MI08, 1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>7300</td>
<td>23.6</td>
<td></td>
<td></td>
<td></td>
<td>$\gamma_1$</td>
<td>(1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>7500</td>
<td>≈ 8000</td>
<td>23.8</td>
<td>(0)</td>
<td>$1^-, 2^-$</td>
<td>$\gamma_1$</td>
<td>(1963MI08, 1963PE15, 1966FI1B, 1968BL1E, 1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>(11100)</td>
<td>(27.0)</td>
<td></td>
<td></td>
<td></td>
<td>$\gamma_1$</td>
<td>(1970FI1B)</td>
<td></td>
</tr>
<tr>
<td>13000</td>
<td>broad</td>
<td>28.6</td>
<td></td>
<td></td>
<td></td>
<td>(1967FE04)</td>
<td></td>
</tr>
</tbody>
</table>

$a$ See (1959AJ76).

$b$ (1964SC19). See however reaction 18.

c $\gamma_0, \gamma_1, \gamma_3, \gamma_4$ represent transitions to $^8$Be*(0, 2.9, 16.6, 16.9), respectively.
Table 8.7: Electromagnetic transitions in $^8$Be $^a$

| Transition | $\Gamma_\gamma$ (eV) | $|M|^2$ (W.u.) | Refs. |
|------------|----------------------|----------------|-------|
| $17.6 \rightarrow 0$ | 16.7 | 0.15 | (1949FO18, 1961ME10, 1968PA09) |
| $17.6 \rightarrow 2.9$ | $8.15 \pm 0.07$ (M1) | 0.12 | (1961ME10, 1968PA09) |
| $17.6 \rightarrow 16.6$ | $0.032 \pm 0.003$ | 1.48 $\pm$ 0.15 (M1) | (1969SW01) |
| $17.6 \rightarrow 16.9$ | 0.021 $\pm$ 0.004 $^b$ | 1.1 | (1968PA09) |
| $17.6 \rightarrow 16.9$ | $0.0013 \pm 0.0003$ | $0.15 \pm 0.04$ (M1) | (1969SW01) |
| $17.6 \rightarrow 16.9$ | $0.0016 \pm 0.0004$ $^b$ | 0.22 | (1968PA09) |
| $18.15 \rightarrow 0$ | 3.0 | 0.03 | (1970FI1B) |
| $18.15 \rightarrow 2.9$ | 3.8 | 0.05 | (1970FI1B) |
| $18.15 \rightarrow 16.6$ | $0.077 \pm 0.019$ | 1.04 $\pm$ 0.26 (M1) | (1969SW01) |
| $18.15 \rightarrow 16.9$ | $0.084 \pm 0.018$ $^b$ | 1.2 | (1968PA09) |
| $18.15 \rightarrow 16.9$ | $0.062 \pm 0.007$ | $1.51 \pm 0.17$ (M1) | (1969SW01) |
| $18.15 \rightarrow 16.9$ | $0.041 \pm 0.011$ $^b$ | 1.1 | (1968PA09) |
| $18.9 \rightarrow 16.6$ | 0.168 | 0.053 (E1) | (1969SW01) |
| $18.9 \rightarrow 16.9$ | 0.099 | 0.045 (E1) | (1969SW01) |
| $19.06 \rightarrow 2.9$ | 10.5 | 0.12 | (1970FI1B) |

$^a$ See also (1966LA04).

$^b$ Values listed by (1968PA09) multiplied by factor 0.56: see (1969SW01), p. 1019.

For a review of the earlier work, see (1959AJ76, 1966LA04). See also (1973SU1E), (1966EV1B, 1966PE1D, 1966WA1C, 1969KA1J), and (1973AS02; theor.).

17. $^7$Li(p, n)$^7$Be

$Q_m = -1.64422 \quad E_b = 17.256$

Recent measurements of cross sections have been made for $E_p = 1.9$ to 2.36 MeV (1967BE61; $\sigma_1$), 1.93 to 2.66 MeV (1969LE23), 2.1 to 3.8 MeV (1971BU1D), 2.4 to 6.0 MeV (1972PR03; $n_1\gamma$), 2.6 to 5.4 MeV (1972EL19; $n_0$), 3 to 10 MeV (1966HA1J; $n_1\gamma$), 3.2 to 5.4 MeV (1972EL19; $n_1$), 23 to 52 MeV (1967LO07; $n_1\gamma$), and 30 and 50 MeV (1969CL06; $n_0$, $n_1$). See also (1973WA34). The shape of the neutron yield near threshold has been studied by (1966PA03). Polarization measurements are reported at $E_p = 2.05$ to 3.00 MeV (1973RO35, 1973RO2E), 3.0 to 4.0 MeV (1971HA27; $n_0$, $n_1$), 3.0 to 5.5 MeV (1971TH07; $n_0$, $n_1$) and 30 and 50 MeV (1969RO20). For a report on the earlier yield and polarization measurements, see (1966LA04). For angular distributions, see $^7$Be.

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Table 8.8: $^8\text{Be}$ levels from $^7\text{Li}(p, n)^7\text{Be}$

<table>
<thead>
<tr>
<th>$E_p$ (MeV)</th>
<th>$\Gamma_{\text{lab}}$ (keV)</th>
<th>$^8\text{Be}^*$</th>
<th>$J^\pi$</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.88</td>
<td>55 ± 20</td>
<td>18.90</td>
<td>2$^-$</td>
<td>(1974AR10) $^a$</td>
</tr>
<tr>
<td>2.25</td>
<td>220</td>
<td>19.22</td>
<td>3$^+$</td>
<td>(1957NE1A, 1961BE05)</td>
</tr>
<tr>
<td>2.6 $^b$</td>
<td>≈ 750</td>
<td>19.5</td>
<td>1$^-$</td>
<td>(1972PR03) $^a$</td>
</tr>
<tr>
<td>3.0</td>
<td>≈ 1250</td>
<td>19.9</td>
<td>(2$^+$)</td>
<td>(1972PR03) $^a$</td>
</tr>
<tr>
<td>4.9</td>
<td>1100</td>
<td>21.5</td>
<td>3(+)</td>
<td>(1959GI47, 1963BO06)</td>
</tr>
<tr>
<td>5.5</td>
<td>broad</td>
<td>22.1</td>
<td>c</td>
<td>(1972PR03)</td>
</tr>
</tbody>
</table>

$^a$ See also (1966LA04).

$^b$ $\gamma_{n1}^2$ and $\gamma_{p1}^2 \approx 1\%$ of the Wigner limit (1972PR03).

$^c$ The broad dip in the $n_1$ yield at the same energy as the broad bump in the $p_1$ yield may be due to interference of two $2^+$ states (1972PR03).

The yield of ground state neutrons ($n_0$) rises steeply from threshold and shows pronounced resonances at $E_p = 2.25$ and 4.9 MeV (1963BO06). The yield of $n_1$ also rises steeply from threshold (1964BU08) and exhibits a broad maximum near $E_p = 3.2$ MeV (1961BE05, 1972PR03) and a broad dip at $E_p \approx 5.5$ MeV, also observed in the $p_1$ yield (1972PR03).

Multi-channel scattering length approximation analysis of the $2^-$ partial wave near the $n_0$ threshold indicates that the $2^-$ state at $E_x = 18.9$ MeV is virtual relative to the threshold and that its width $\Gamma = 50 \pm 20$ keV (1974AR10). The ratio of the cross section for $^7\text{Li}(p, \gamma)^8\text{Be}^*(18.9)$ to the thermal neutron capture cross section $^7\text{Be}(n, \gamma)^8\text{Be}^*(16.6 + 16.9)$ [obtained by (1969SW01)] provides a rough estimate of the isospin impurity of $^8\text{Be}^*(18.9)$: $\sigma_{p,\gamma}/\sigma_{n,\gamma} \approx 1.5 \times 10^{-5}$ and therefore the $T = 1$ isospin impurity is $< 4\%$ in intensity (1974AR10).

The structure at $E_p = 2.25$ MeV is ascribed to a $3^+$, $T = (1)$, $l = 1$ resonance with $\Gamma_n \approx \Gamma_p$ and $\gamma_n^2/\gamma_p^2 = 3$ to 10: see (1966LA04). See also (1973RO35). At higher energies the broad peak in the $n_0$ yield at $E_p = 4.9$ MeV can be fitted by $J^\pi = 3(+)$ with $\Gamma = 1.1$ MeV, $\gamma_n^2 \approx \gamma_p^2$ (1963BO06). The behavior of the $n_1$ cross section can be fitted by assuming a $1^-$ state at $E_x = 19.5$ MeV and a $J = 0, 1, 2$, positive-parity state at 19.9 MeV [presumably the 20.2 MeV state reported in reaction 4]. In addition the broad dip at $E_p \approx 5.5$ MeV may be accounted for by the interference of two $2^+$ states (1972PR03). See Table 8.8.

The ratio of the cross sections of the $(p, n_1)$ reaction to $^7\text{Be}^*(0.43)$ to that for the $(p, p_1)$ reaction to the analog state $^7\text{Li}^*(0.48)$ has been measured for $E_p = 2.4$ to 6.0 MeV (1972PR03), 3 to 10 MeV (1966HA11) and 23 to 52 MeV (1967LO07). At the lower energies it deviates markedly from unity and varies strongly with energy (1966HA11). At the higher energies the measurements seem to indicate that the spin-flip, isospin-flip part of the effective interaction is essentially independent of energy while the pure central part appears to decrease as the energy increases (1967LO07).
also (1969CL06).


18. (a) $^7$Li(p, p)$^7$Li

$$E_b = 17.256$$

(b) $^7$Li(p, p')$^7$Li*

Absolute differential cross sections for elastic scattering have been reported for $E_p = 0.4$ to 12 MeV (1953WA27, 1956MA12, 1965GL03), 14.5, 20.0 and 31.5 MeV (1956KI54) and more recently at 0.85 to 2.0 MeV (1966BA1Q), at 1.36 MeV (1969LE08) and at 6.868 MeV (1971BI11). The yield of inelastically scattered protons (p$_1$, to $^7$Li*(0.48)) and of 0.48 MeV $\gamma$-rays have been measured in the range $E_p = 0.8$ to 12 MeV (1953WA27, 1956MA12, 1965GL03), 14.5, 20.0 and 31.5 MeV (1956KI54) and more recently at 0.85 to 2.0 MeV (1966BA1Q, 1970RO22, 1971BI11; p$_0$).

Polarization measurements are reported at $E_p = 0.67$ to 2.45 MeV (1973BR13; p$_0$), 2.7 to 10.6 MeV (1969KI04; p$_0$, p$_1$), 14.5 MeV (1965RO22; p$_0$), 49.8 MeV (1971MA13, 1971MA44; p$_0$, p$_1$), 152 MeV (1966RO1C; p$_0$) and 155 MeV (1968GE04; p$_0$, p$_2$). For earlier measurements see (1966LA04). For a summary of angular distribution studies see $^7$Li.

Anomalies in the elastic scattering appear at $E_p = 0.44$, 1.03, 1.88, 2.1, 2.5, 4.2 and 5.6 MeV. Resonances at $E_p = 1.03$, 3 and 5.5 MeV and an anomaly at $E_p = 1.88$ MeV appear in the inelastic channel. A phase shift analysis and a review of the existing cross section data by (1973BR13) show that the 0.44 and 1.03 MeV resonances are due to $1^+$ states which are a mixture of $^5P_1$ and $^3P_1$ with a mixing parameter of $+25^\circ$; that the $2^-$ state at the neutron threshold ($E_p = 1.88$ MeV) has a width of about 50 keV [see also reaction 17]; and that the $E_p = 2.05$ MeV resonance corresponds to a $3^+$ state. The anomalous behavior of the $^5P_3$ phase around $E_p = 2.2$ MeV appears to result from the coupling of the two $3^+$ states [resonances at $E_p = 2.05$ and 2.25 MeV]. The $^3S_1$ phase begins to turn positive after 2.2 MeV suggesting a $1^-$ state at $E_p = 2.5$ MeV (1973BR13): see Table 8.9 and (1972PR03).

An attempt has been made to find the $T = 2$ analog of the ground state of $^8$He: no resonances were observed in either the p$_0$ or the p$_1$ yield for $E_p = 11.1$ to 11.9 MeV (1968HA1H). Measurements of the intensity ratios of the reactions (p, p$_1$) and (p, n$_1$) have been made by (1966HA1J, 1967LO07, 1972PR03): see reaction 17.

See also (1967CA1G, 1972RU1C) and (1967JO01, 1967SA1C, 1969WA11; theor.).

19. $^7$Li(p, d)$^6$Li

$$Q_m = -5.026$$

$$E_b = 17.256$$

See (1969KO1P; theor.) and $^6$Li.

20. $^7$Li(p, t)$^6$Li

$$Q_m = -4.43$$

$$E_b = 17.256$$
### Table 8.9: $^8$Be levels from $^7$Li(p, p)$^7$Li and $^7$Li(p, p)$^7$Li*

<table>
<thead>
<tr>
<th>$E_p$ (MeV)</th>
<th>$\Gamma_{\text{lab}}$ (keV)</th>
<th>$^8$Be*</th>
<th>$J^\pi$</th>
<th>$\Gamma_{p'}$ (keV)</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.441</td>
<td>12.2 b</td>
<td>17.642</td>
<td>1+</td>
<td></td>
<td>(1953CH1A, 1953WA27, 1973BR13)</td>
</tr>
<tr>
<td>1.030 ± 0.005</td>
<td>168</td>
<td>18.157</td>
<td>1+</td>
<td>6</td>
<td>(1954MO04, 1955LI1B, 1973BR13)</td>
</tr>
<tr>
<td>2.05</td>
<td>≈ 400</td>
<td>19.05</td>
<td>3+</td>
<td>small</td>
<td>(1956MA12, 1957NE1A, 1973BR13)</td>
</tr>
<tr>
<td>2.25</td>
<td></td>
<td>19.22</td>
<td>3+</td>
<td>small</td>
<td>(1956MA12, 1957NE1A, 1973BR13)</td>
</tr>
<tr>
<td>2.5 c</td>
<td>≈ 750</td>
<td>19.4</td>
<td>1−</td>
<td>res.</td>
<td>(1972PR03, 1973BR13)</td>
</tr>
<tr>
<td>d</td>
<td>1800 ± 200</td>
<td>20.9 e</td>
<td>4−</td>
<td>(res.)</td>
<td>(1965GL03)</td>
</tr>
<tr>
<td>4.2 ± 0.2</td>
<td>broad</td>
<td>22.2</td>
<td>f</td>
<td>res.</td>
<td>(1965GL03, 1972PR03)</td>
</tr>
<tr>
<td>5.6</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

a (p, n) threshold: see reaction 17.
b $\theta^2_p = 0.064$.
c See also Table 8.8, $\gamma^2_{\alpha n}$, and $\gamma^2_{p1}$ ≈ 1% of the Wigner limit (1972PR03).
d A 2+ state at $E_x \approx 20$ MeV appears to be necessary to account for the cross sections: see Table 8.3 and reaction 4 (1972PR03).
e Reduced width is 70% of the Wigner limit (1965GL03).
f May be due to two 2+ states (1972PR03). See also reaction 17.

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See $^5$Li.

21. $^7$Li(p, $\alpha$)$^4$He

$Q_m = 17.348$  
$E_b = 17.256$

The cross section follows the expression $E^{-1}e^{-B/\sqrt{E}}$, with $B = 91.5 \pm 4.5$ keV$^{1/2}$, in the range $E_p = 23$ to 50 keV. The cross section in that interval rises from 0.013 to 2.4 µb (1967FI05). In the range $E_p = 131$ to 561 keV, the cross section increases from 0.16 ± 0.02 to 3.7 ± 0.4 mb (1971SP05; and T.A. Tombrello, private communication). The cross section has also been measured for $E_p = 0.04$ to 0.13 MeV by (1969LE22). Taking into account $^8$Be $J^\pi = 2^+$ levels at 16.7, 16.9 and 20.6 MeV, (1972BA41) has made an $R$-matrix fit to the revised data of (1971SP05) and has obtained a quadratic energy dependence for the $S$-factor: $S = 0.065[1 + 1.82E - 2.51E^2]$ MeV · b, over the energy range $E_p = 0$ to 600 keV.

Excitation functions and angular distributions have been measured at many energies up to 18.6 MeV; see (1966LA04) for earlier references. Recently, differential cross-section measurements are reported by (1966MA03, 1969SW01: $E_p = 0.4$ to 2.45 MeV), (1969LE08: $E_p = 1.36$ MeV).
and (1967CR05; \( E_p = 41.3 \) and 45.2 MeV). Polarization measurements have been carried out for 
\( E_p = 0.8 \) to 3 MeV (1968PE03), 2.7 to 10.6 MeV (1969KI04), 3.00 to 10.04 MeV (1968PL01), 5.5 
to 6.7 MeV (1966BO09), and 7.4 to 10.4 MeV (1968AR04). See also (1966LA04) for a listing of 
the earlier references and (1966DA1B, 1971PL1C). In the range \( E_p = 3 \) to 10 MeV the asymmetry 
has one broad peak in the angular distribution at all energies except near 5 MeV; the peak value is 
0.98 ± 0.04 at 6 MeV and is essentially 1.0 for \( E_p = 8.5 \) to 10 MeV (1968PL01, 1969KI04) [see 
Fig. 12 in (1969KI04) and Fig. 6 (1968PL01) for contour maps of the asymmetry].

Broad resonances are reported to occur at \( E_p = 3.0 \) MeV, \( \Gamma \approx 1 \) MeV (1948HE1A) and at 
\( E_p = 5.6 \) MeV, \( \Gamma \approx 1 \) MeV (1961HA27, 1962TE04, 1964MA51). Some structure is also reported 
near \( E_p = 6.0 \) to 6.5 MeV, and at \( E_p = 9.0 \) MeV (1964MA51). The latter is also reflected in the 
behavior of the \( A_2 \) coefficient (1968PL01). The experimental data on yields and on polarization 
have been analyzed by (1970KU1H, 1971KU10): the data appear to require including two \( 0^+ \) states 
[at \( E_x \approx 19.7 \) and 21.8 MeV] with very small \( \alpha \)-particle widths, and four \( 2^+ \) states [at \( E_x \approx 15.9, 
20.1, 22.2 \) and 25 MeV]. See, however, reaction 4 and (1972BA83). A \( 4^+ \) state near 20 MeV was 
also introduced in the calculation but its contribution was negligible. The observed discrepancies 
are said to be probably due to the assumption of pure \( T = 0 \) for these states (1971KU10).


At \( E_p = 9.1 \) MeV, \( \alpha \)-particle spectra are discussed in terms of the first excited state of \(^4\)He 

22. (a) \(^7\)Li(d, n)\(^8\)Be \quad Q_m = 15.031
(b) \(^7\)Li(d, n)\(^4\)He\(^\alpha\) \quad Q_m = 15.1233

At \( E_d = 2 \) MeV, recoil proton spectra show only the ground state and \(^8\)Be*(2.9). No other 
groups with \( E_x < 9 \) MeV appear with intensity > 10% of \( n_0 \). The spectrum yields \( E_x = 3.1 \pm 0.1, 
\Gamma = 1.75 \pm 0.1 \) MeV (1964JO04) [(1971RO05) report \( E_x = 3.10 \pm 0.09, \Gamma = 1.74 \pm 0.08 \) MeV]. See 
Table 8.4. At higher deuteron energies the population of \(^8\)Be*(16.6, 16.9, 17.6, 18.2, 18.9, 19.1, 
19.2) is reported and \( l_p = 1 \) is obtained for the transitions to \(^8\)Be*(16.6, 17.6, 18.2): see (1960DI02, 
1966DI1B, 1967KE1A, 1967KE1F). Angular distributions of the \( n_0 \) and \( n_1 \) groups to \(^8\)Be*(0, 2.9) 
are reported by (1966JU1A: \( E_d = 0.7 \) and 0.8 MeV), (1969NU1C: \( E_d = 0.90 \) to 1.09 MeV), 
(1966MI09: \( E_d = 1 \) MeV), and by (1970OSZY: \( E_d = 1.62 \) to 2.97 MeV). See also (1966LA04) 

Reaction (b) appears to proceed primarily by sequential decay via \(^8\)Be*(2.9, 16.6, 16.9) and 
1973MC13). However, (1969HO11) deduce the involvement of a state with \( E_x = 11.4 \pm 0.05 \) MeV, 
\( \Gamma_{c.m.} = 2.8 \pm 0.2 \) MeV. See also (1973KA32). Attempts to observe \( n-\alpha \) rescattering ("proximity 
scattering") proceeding via \(^8\)Be*(16.6, 16.9) have been unsuccessful: see (1968VA12, 1971SW10, 
1972BR08). See also (1969TH02, 1971TH08) and the discussion in \(^5\)He.

23. (a) $^7$Li($^3$He, d)$^8$Be $Q_m = 11.762$
(b) $^7$Li($^3$He, pn)$^8$Be $Q_m = 9.538$

Deuteron groups are observed to $^8$Be*(0, 2.9, 16.6, 16.9, 17.6). The group to $^8$Be*(2.9) is well fitted by $E_x = 2.82$ MeV, $\Gamma = 1.27$ MeV (1971PI06). See also reaction 13 (1969NU01). The $J^\pi = 1^+$ mixed isospin state have $E_x = 16.627 \pm 0.005$ and $16.901 \pm 0.005$ MeV and $\Gamma = 113 \pm 3$ and $77 \pm 3$ keV (1967MA12): see also Table 8.5 and (1971PI06). Angular distributions have been measured at $E(3\text{He}) = 0.90$ and 1.10 MeV (1971ST35; d$_0$, d$_1$), 3 MeV (1972LI31; d$_0$, d$_1$), and at 10 MeV by (1970DI12, 1970DI1F; d$_0$, d$_1$) and by (1968CO07; d to $^8$Be*(16.6, 16.9)). Spin-dependent effects in the angular distributions of d$_0$ and d$_1$ obtained by (1963WE1B) at 24.3 MeV are discussed by (1967SI1A). The angular distribution to $^8$Be*(16.6) is forward peaked, that for $^8$Be*(16.9) is roughly isotropic (1968CO07). See also (1964MA57).

The decay of various $^9$Be states to $^8$Be*(0, 2.9) has been studied by (1966CH20, 1968CO08, 1972MC1E): see $^9$Be and Table 9.5. See also (1968LI1D, 1970LI1Q) and (1967CO1L; theor).

24. (a) $^7$Li($\alpha$, t)$^8$Be $Q_m = -2.559$
(b) $^7$Li($\alpha$, t$\alpha$)$^4$He $Q_m = -2.4668$

The angular distributions of the t$_0$ group have been measured at many energies up to 48 MeV: see (1966LA04). Recently, measurements have been carried out at $E_{\alpha} = 23.2$ and 25.0 MeV (1973VA1A; t$_0$, t$_1$), 30 MeV (1972ME07; t$_0$, t$_1$; PWBA and DWBA analysis) and at 50 MeV (1970LA14; t$_0$). The ground state of $^8$Be does decay isotropically in the c.m. system and therefore $J^\pi = 0^+$ (1970LA14). Spin-dependent effects in the angular distributions for t$_0$ and t$_1$ obtained by (1963WE1B: $E_{\alpha} = 28$ MeV) are discussed by (1967SI1A). At $E_{\alpha} = 10$ MeV an anomaly (“ghost”) is observed in the $^8$Be excitation spectrum at $E_x \approx 0.5$ MeV. It may be due to interference of the 0$^+$ states $^8$Be*(0, 6.) [see reaction 4] or to thresholds of particle channels (1971BE52).

In reaction (b), sequential decay is observed at $E_{\alpha} = 50$ MeV, via $^8$Be*(0, 2.9, 11.4, 16.6, 16.9, 19.9) (1970LA14). See also (1968BE1Q, 1968MA25).

25. $^7$Li($^7$Li, $^6$He)$^8$Be $Q_m = 7.278$

At $E(7\text{Li}) = 1.4, 1.7$ and 1.8 MeV, the angular distributions of $^6$He ions leaving $^8$Be in its ground and 2.9 MeV states are essentially isotropic (1968ST12). See also (1966LA04).
26. (a) $^7\text{Be}(n, p)^7\text{Li}$
   \[ Q_m = 1.64422 \quad E_h = 18.900 \]
(b) $^7\text{Be}(n, \alpha)^4\text{He}$
   \[ Q_m = 18.992 \]
(c) $^7\text{Be}(n, \gamma\alpha)^4\text{He}$
   \[ Q_m = 18.992 \]

At thermal energies, the $(n, p)$ cross section is $(4.8 \pm 0.9) \times 10^4$ b \((1955\text{HA34}, 1973\text{MU14})\), the $(n, \alpha)$ cross section is $\leq 0.1$ mb \((1962\text{BA1B}, 1963\text{BA34})\) and the $(n, \gamma\alpha)$ cross section is 155 mb \((1963\text{BA34})\). These values, and comparison of the $(p, n)$ cross section with that of reaction (a), support the $J^\pi = \frac{3}{2}^-$ assignment for $^7\text{Be}(0)$ \((1957\text{NE1A}, 1963\text{BA34})\). The role of these reactions in astrophysical phenomena is discussed by \((1968\text{FO1A}, 1969\text{BA1N})\). See also \((1959\text{AJ76})\) and reaction 17.

27. $^7\text{Be}(d, p)^8\text{Be}$
   \[ Q_m = 16.676 \]

For $E_d = 0.8$ to 1.7 MeV, proton groups are observing corresponding to the ground state and $^8\text{Be}^*(2.9)$: derived parameters for the latter are shown in Table 8.4 \((1959\text{SP1A}, 1960\text{KA17})\). See also \((1969\text{BA43})\) and $^8\text{B}^*(\beta^+)$. Studies of the distribution of recoil momenta and neutrino recoil correlation indicate that the decay is overwhelmingly GT, axial vector [see reaction 1 in $^8\text{Li}$] and that the ground state of $^8\text{Li}$ has $J^\pi = 2^+$: see \((1966\text{LA04})\).

Angular correlations have been measured for the decays of $^8\text{Li}$ and $^8\text{B}$ as a test of the conserved vector current theory of $\beta$-decay. The values of the coefficients are displayed in Table 8.10. See also \((1973\text{TR1J}, 1973\text{TR1K}, 1973\text{TR1L})\). The experimental value of $\delta$ \([\delta \equiv B(^8\text{Li}) - B(^8\text{B})]\) is $(5.4 \pm 0.4) W_\beta$, consistent with CVC theory \((1966\text{EI02})\).

A recent asymmetry measurement is reported by \((1971\text{VA19})\). See also \((1971\text{VA1E}, 1973\text{NE10})\). Measurements of the excitation spectra in the decays of $^8\text{Li}$ and $^8\text{B}$ show no evidence for second class currents: $|g_{TT}| < 7 \times 10^{-4}$ \((1971\text{WI05})\). See also \((1966\text{JA1C})\) and \((1960\text{KU05}, 1960\text{WE1A}, 1966\text{BA26}, 1966\text{LI1C}, 1968\text{KR10}, 1969\text{BA43}, 1970\text{DA21}, 1971\text{LI1H}, 1971\text{WI18}, 1971\text{WI1C}, 1972\text{EM02}, 1972\text{HO23}, 1972\text{MI1M}, 1972\text{WI28}, 1972\text{WI1C}, 1973\text{EM1B}, 1973\text{HA49}, 1973\text{TO14}, 1973\text{WI11}, 1974\text{WI1L}; \text{theor.})\).
Table 8.10: $\alpha$-$\beta$ angular correlation coefficients in $^8\text{Li}, ^8\text{B}$

<table>
<thead>
<tr>
<th>Nuclide</th>
<th>$A/W_\beta$</th>
<th>$B/W_\beta$</th>
<th>$W_\beta$ (MeV)</th>
<th>$\delta/W_\beta$</th>
<th>Refs.</th>
</tr>
</thead>
<tbody>
<tr>
<td>$^8\text{Li}$</td>
<td>$(-8.7 \pm 0.7) \times 10^{-3}$</td>
<td>$(5.7^{+2.9}_{-1.9}) \times 10^{-3}$</td>
<td>7.0</td>
<td></td>
<td>(1960KR03)</td>
</tr>
<tr>
<td>$^8\text{Li}$</td>
<td>$(-8.3 \pm 1.1) \times 10^{-3}$</td>
<td>$(+3.2 \pm 0.6) \times 10^{-3}$</td>
<td>11</td>
<td>$(7.0 \pm 1.2) \times 10^{-3}$</td>
<td>(1962NO02)</td>
</tr>
<tr>
<td>$^8\text{Li}$</td>
<td>$(-9.7 \pm 0.7) \times 10^{-3}$</td>
<td>$(+3.7 \pm 1.0) \times 10^{-3}$</td>
<td>7.5</td>
<td></td>
<td>(1963GR11)</td>
</tr>
<tr>
<td>$^8\text{B}$</td>
<td>$(-8.7 \pm 0.9) \times 10^{-3}$</td>
<td>$(+3.1 \pm 0.3) \times 10^{-3}$</td>
<td>6.6</td>
<td>$(5.4 \pm 0.4) \times 10^{-3}$</td>
<td>(1966EI02)</td>
</tr>
<tr>
<td>$^8\text{B}$</td>
<td>$(-11.1 \pm 1.3) \times 10^{-3}$</td>
<td>$(-2.3 \pm 0.3) \times 10^{-3}$</td>
<td>7.0</td>
<td></td>
<td>(1966EI02)</td>
</tr>
</tbody>
</table>

$^a W(\theta) = 1 + A \cos \theta + B \cos^2 \theta.$

$^b \delta \equiv B(^8\text{Li}) - B(^8\text{B}).$

29. $^8\text{B}(\beta^+)^9\text{Be}$

$Q_m = 17.981$

The decay proceeds mainly to $^8\text{Be}^*(2.9)$ [see Table 8.4 for its parameters]. Detailed study of the high energy portion of the $\alpha$-spectrum reveals a maximum near $E_\alpha = 8.3$ MeV, corresponding to transitions to $^8\text{Be}^*(16.63)$, for which parameters $E_x = 16.67$ MeV, $\Gamma = 150$ to 190 keV or $E_x = 16.62$ MeV, $\Gamma = 95$ keV are derived. Using $\tau_{1/2} = 769 \pm 4$ msec, $\log ft = 2.9$. The low $ft$ value supports the identification $J^\pi = 2^+$; $T = 1$ for $^8\text{Be}^*(16.63)$ (1964MA35). See, however, (1965MA1G). The energy distribution of $\alpha$-particles has also been measured by (1969CL10). Analysis of this data and of data from $\alpha$-$\alpha$ scattering in a three level $R$-matrix formalism indicate a $2^+$ state of $^8\text{Be}$ at $E_x = 12.0^{+3.0}_{-7.0}$ MeV and of $\Gamma = 14^{+3}_{-4}$ MeV ($a_2 = 6.0 \pm 0.5$ fm) (1969CL10). See also (1969BA43; theor.).

For angular correlation measurements see reaction 28 (1966EI02) and Table 8.10. See also (1973TR1J, 1973TR1K, 1973TR1L).

30. (a) $^9\text{Be}(\gamma, n)^8\text{Be}$

$Q_m = -1.6651$

(b) $^9\text{Be}(n, 2n)^8\text{Be}$

$Q_m = -1.6651$

(c) $^9\text{Be}(p, pn)^8\text{Be}$

$Q_m = -1.6651$

(d) $^9\text{Be}(t, tn)^8\text{Be}$

$Q_m = -1.6651$

(e) $^9\text{Be}(\alpha, \alpha n)^8\text{Be}$

$Q_m = -1.6651$

For reaction (a) see (1966DE07, 1968AD09, 1969GA1M) and (1965BO1B, 1969AU05; theor.). See also $^9\text{Be}$.

Reaction (b) appears to proceed largely via excited states of $^9\text{Be}$, with subsequent decay to $^8\text{Be}$, mainly $^8\text{Be}^*(2.9)$: see (1966LA04), $^9\text{Be}$ and $^{10}\text{Be}$. At $E_n = 14$ MeV the cross section at

For reaction (c) see (1966NO1A; theor.) and ⁹Be. For reaction (d) see (1967SE11). For reaction (e) see (1971GU15, 1973GE1J) and ⁹Be.

31. (a) ⁸Be(p, d)⁶Be \( Q_m = 0.5595 \)
(b) ⁸Be(p, d)⁴He⁴He \( Q_m = 0.651 \)

\( Q_0 = 559.0 \pm 1.1 \text{ keV (1966RE02, 1967ST30); } \)
\( Q_0 = 559.6 \pm 0.6 \text{ keV (1967OD01); see also (1967SP09).} \)

Angular distributions of deuteron groups have been reported at \( E_p = 0.11 \) to 0.55 MeV (1973SI27; \( d_0 \)), 0.30 to 0.90 MeV (1968BE1N; \( d_0 \)), 5 to 11 MeV (1972HU03; \( d_0 \)) [analysis by DWBA and BHMM (1967BU23); derived spectroscopic factors], 13.0, 14.0, 15.0 and 21.35 MeV (1972VO1H; \( d_0 \)), 17.0, 21.0, 25.0, 29.1 MeV (1973MO01; \( d_0 \), \( d_1 \)), 33.6 MeV (1967KU10; 1970KU1D: deuterons to ⁸Be*(0, 2.9, 16.95, 17.62, 18.18, 19.21); also derived spectroscopic factors) [also saw ⁸Be*(11.4)]; determined \( \Gamma_{e.m.}(16.95) = 103 \pm 15 \text{ keV, } \Gamma_{e.m.}(19.21) = 208 \pm 30 \text{ keV}, \) 40.8 MeV (1966MA22: deuterons to ³Be*(16.63, 16.91)), 46 MeV (1967VE01: deuterons to ³Be*(0, 2.9, 16.9, 17.6, 18.2, 19.1) [also report ³Be*(24.5)], 100 MeV (1968LE01: deuterons to ³Be*(0, 2.9, 16.9, 18.9)) [also saw ³Be*(11.0, 23.0, 26.0)], 155 MeV (1969BA05, 1969TO1A: deuterons to ³Be*(0, 2.9) [\( \Gamma = 2 \pm 0.1 \text{ MeV}, \) 11.5 [\( \Gamma = 8 \pm 1 \text{ MeV}, \) 16.8 \( \pm 0.2, \) 18.9 \( \pm 0.3 \) [also saw ³Be*(17.6 \( \pm 0.4, 21.5 \pm 0.3 \)]] and 185 MeV (1969SU02: deuterons to ³Be*(0, 2.94 \( \pm 0.08, \) 11.3 \( \pm 0.3, \) 16.87 \( \pm 0.06, \) 17.58 \( \pm 0.08, \) 18.10 \( \pm 0.10, \) 19.16 \( \pm 0.07, \) 22.0 \( \pm 0.15, \) 22.9 \( \pm 0.15 \)]] [also report ³Be*(20.0 \( \pm 0.2 \) (?)); \( \Gamma_{e.m.}(2.9) = 1.5 \pm 0.1 \text{ MeV; } \Gamma_{e.m.}(22.0, 22.9) \geq 1 \text{ MeV (the angular distributions for these two states are not clear cut)].} \) (1971SC26) have analyzed the angular distributions obtained by (1967VE01, 1968LE01) using DWBA with a local-energy approximation and have derived spectroscopic factors. With the exception of ³Be*(11.4, 22.0, 22.9) the angular distributions are consistent with \( l = 1 \). The yield of the deuterons corresponding to ³Be*(16.63) is very low: [\( \approx 5\% \) compared to ³Be*(16.91)] as expected by predictions of the cluster model (1966MA22: \( E_p = 40.8 \) MeV). See also (1967KU10) and reaction 21 in ⁹Be in (1966LA04).

Anomalies in the deuteron spectrum between the \( d_0 \) and the \( d_1 \) groups have been reported at various energies [see (1966LA04) and (1967FI1D, 1967HA1K, 1971BE52, 1971MI1C)]. The shape of the deuteron spectrum near ³Be*(2.9) requires \( a_2 \approx 7.1 \text{ fm (1969BA43).} \) See also reaction 4 and (1968BA2D). At \( E_p = 17 \text{ MeV, for the transitions to ³Be*(0, 2.9) the ratios of } \sigma(p,d)/\sigma(p,d) = 11.8 \text{ and } 14.1, \text{ respectively (1967CO09, 1969CO06).} \) See also (1968NA1A; theor.).

Reaction (b) at \( E_p = 9 \text{ MeV is dominated by strong final state interactions through ³Be*(0, 2.9) and ⁴Li*(2.19) with little or no yield from a direct three-body decay (1971EM01).} \) See also (1967FI1D). See also ¹⁰B and (1966CA1E, 1966LA20, 1967OG1A, 1967RO07, 1968TI1A, 1972QU01) and (1967BA1M, 1967JO1D, 1968BO1P, 1970BO1K; theor.).
32. (a) $^9$Be(d, t)$^8$Be \[ Q_m = 4.5925 \]
(b) $^9$Be(d, t)$^4$He$^4$He \[ Q_m = 4.684 \]
\[ Q_0 = 4591.7 \pm 3.1 \text{ keV} \] (1967OD01; see also (1967SP09)).

At $E_d = 11.8$ MeV, angular distributions have been obtained for the tritons to $^8$Be$^*(0, 2.9)$ (1967FI07): $S = 0.51$ and 0.75, respectively (DWBA analysis). At $E_d = 38$ MeV, angular distributions of the tritons to $^8$Be$^*(16.91, 17.64, 19.0)$ have been compared with those of the $^3$He to the analog states in $^8$Li. The cross-section ratios $\sigma_{17.64}/\sigma_{0.98}^{(8\text{Li})} = 0.45 \pm 0.04$ and $\sigma_{16.91}/\sigma_{0}^{(8\text{Li})} = 0.75 \pm 0.04$, consistent with the pure $T = 1$ nature of $^8$Be$^*(17.64)$ and the mixed $T$ nature of $^8$Be$^*(16.91)$ [and $^8$Be$^*(16.63)$] (1966GA21; abstract). Angular distributions have also been measured at $E_d = 0.3$ to 1.0 MeV (1968BE1E; $t_0$), 0.9 to 3.1 MeV (1973SA1Q; $t_0$), 15.0 MeV (1969AR11B; $t_0$) and at many other energies up to $E_d = 20$ MeV; see (1966LA04). The ghost anomaly which is seen near the $t_0$ group has been studied at $E_d = 2.5$ MeV: it is interpreted as being due to an extreme threshold effect (1971BE52). See also (1967DE1J) and (1970BO1K, 1973HE1J; theor.). A kinematically complete study of reaction (b) at $E_d = 26.3$ MeV indicates the involvement of $^8$Be$^*(0, 2.9, 11.4, 16.9, 19.9 + 20.1)$. Parameters obtained for $^8$Be$^*(2.9, 11.4)$ are $E_x = 3.20 \pm 0.03$ and 11.70 $\pm 0.07$ MeV, $\Gamma = 1.72 \pm 0.09$ and 4.41 $\pm 0.5$ MeV (1973SO08).

33. (a) $^9$Be($^3$He, $\alpha$)$^8$Be \[ Q_m = 18.9134 \]
(b) $^9$Be($^3$He, $\alpha$)$^4$He$^4$He \[ Q_m = 19.005 \]
(c) $^9$Be($^3$He, $\alpha$p)$^7$Li \[ Q_m = 1.657 \]

Angular distributions have been measured at $E(^3\text{He}) = 3.0$ MeV (1968MO05; $\alpha_{16.91}$), 3.0 and 4.0 MeV (1963DO08; $\alpha_0$, $\alpha_{2.9}$, $\alpha_{16.6}$, $\alpha_{16.9}$, $\alpha_{17.6}$), 18.0, 22.7, 26.7, 32.3 MeV (1965AR07; $\alpha_0$, $\alpha_1$) and 26.7 MeV (1968AR12; $\alpha_{16.9}$, $\alpha_{17.6}$, $\alpha_{18.1}$, $\alpha_{19.2}$). See also (1959AJ76) and (1967SI1A). The parameters of the observed states are shown in Tables 8.4 and 8.5 (1961ER01, 1963DO08).

Reaction (b) has been studied at $E(^3\text{He}) = 1.6$ MeV (1970EH11A), 3.0 MeV (1966SU04, 1968MO05) and 3.0 and 4.0 MeV (1972TA04). See also (1966LA04) and (1967ST1E). The reaction proceeds by sequential decay via $^8$Be$^*(0, 2.9, 11.4, 16.6, 16.9, 19.9, 22.5)$ (1972TA04). The angular correlation via $^8$Be$^*(16.91)$ is consistent with $J^\pi = 2^+$ for that state (1968MO05). $J^\pi = 2^+$ is also indicated for $^8$Be$^*(16.63)$ (1966SU04). See also (1968TH1G). For reaction (c) see (1967ST1D). See also (1964MA57, 1966CA08, 1966DI1C, 1967OG1A, 1970CA28, 1971TR1B, 1972RO1N), (1967HO1C) and (1970BO1K, 1971OS05, 1972TH04, 1973RO28; theor.).

34. (a) $^9$Be($^6$Li, $^7$Li)$^8$Be \[ Q_m = 5.586 \]
(b) $^9$Be($^7$Li, $^8$Li)$^8$Be \[ Q_m = 0.368 \]
(c) $^9$Be($^{12}$C, $^{13}$C)$^8$Be \[ Q_m = 3.281 \]
(d) $^9$Be($^{16}$O, $^{17}$O)$^8$Be \[ Q_m = 2.477 \]
(e) $^9\text{Be}(^{18}\text{O}, \, ^{19}\text{O})^8\text{Be}$ \quad $Q_m = 2.292$

(f) $^9\text{Be}(^{19}\text{F}, \, ^{20}\text{F})^8\text{Be}$ \quad $Q_m = 4.936$

At $E(^6\text{Li}) = 3.5$ MeV the population of $^8\text{Be}^*(2.9)$ is very small but $^8\text{Be}_{\text{g.s.}}$ is involved (1968JA08). See also (1966LA04). For reaction (b) see (1966LE10) and (1968TO1C; theor.). For reaction (c) see (1970BA1J, 1970BA1Y). For reaction (d) see (1968KN1A, 1970BA1J, 1970BA1Y). Reaction (e) has been studied at $E(^{18}\text{O}) = 16$ and 20 MeV (1971KN05). See also (1968FA04). For reaction (f) see (1968FA04).

35. $^{10}\text{B}(\gamma, \, d)^8\text{Be}$ \quad $Q_m = -6.0258$

See $^{10}\text{B}$ and (1959AJ76).

36. (a) $^{10}\text{B}(n, \, t)^8\text{Be}$ \quad $Q_m = 0.2318$

(b) $^{10}\text{B}(n, \, t)^4\text{He}^4\text{He}$ \quad $Q_m = 0.3237$

Angular distributions have been measured at $E_n = 14.4$ MeV (1964VA14; $t_0, \, t_1$). Reaction (b) has been studied at the same energy by (1967VA12). See also (1971MI1H) and (1967BA1E; theor.). See also $^{11}\text{B}$ in (1975AJ02).

37. (a) $^{10}\text{B}(p, \, ^3\text{He})^8\text{Be}$ \quad $Q_m = -0.5320$

(b) $^{10}\text{B}(p, \, \text{pd})^8\text{Be}$ \quad $Q_m = -6.0258$

At $E_p = 49.5$ MeV angular distribution measurements have been carried out for the $^3\text{He}$ groups to $^8\text{Be}^*(0, \, 2.9, \, 16.6, \, 16.9)$: the ratio $d\sigma(16.63)/d\sigma(16.91)$ has a mean value of $0.65 \pm 0.05$ for $\theta = 15^\circ$ to $30^\circ$, suggesting possibly a preferential excitation of the $T = 1$ components of these two states. The ratio of the differential cross sections $d\sigma(p, \, t)$ [to $^8\text{B}_{\text{g.s.}}$] to $d\sigma(p, \, ^3\text{He})$ [to $^8\text{Be}^*(16.63, \, 16.91)$] ($15^\circ$ to $30^\circ$) seems to also suggest this (1971SQ01).

38. (a) $^{10}\text{B}(d, \, \alpha)^8\text{Be}$ \quad $Q_m = 17.822$

(b) $^{10}\text{B}(d, \, \alpha)^4\text{He}^4\text{He}$ \quad $Q_m = 17.9138$

$Q_0 = 17.8186 \pm 0.0041$ (1967OD01).

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Angular distributions have been reported at $E_d = 0.5$ to 1.0 MeV (1968FR07; $\alpha_0$, $\alpha_1$), 0.8 to 2.5 MeV (1968CO31; $\alpha_0$, $\alpha_1$), 3.0 to 7.2 MeV (1967LE1C; $\alpha_0$) and 7.5 MeV (1966BR08; $\alpha$ to $^8\text{Be}^*(16.63, 16.91, 17.64, 18.15)$). At $E_d = 7.5$ MeV the total cross section for formation of $^8\text{Be}^*(16.63)$, $\sigma_t(16.63)$, is about 1.15 $\sigma_t(16.91)$, consistent with the mixed isospin character of these two states. $\sigma_t(18.15)$ is $\approx 0.85 \sigma_t(16.91)$, but the other nearby $1^+$ state $^8\text{Be}^*(17.64)$ has $\sigma_t(17.64) \approx 0.07 \sigma_t(16.91)$, consistent with the nearly pure $T = 1$ nature of $^8\text{Be}^*(17.64)$ (1966BR08). These four states [$^8\text{Be}^*(16.63, 16.91, 17.64, 18.15)$] have been studied for $E_d = 4.0$ to 12.0 MeV. Interference between the $2^+$ states [$^8\text{Be}^*(16.63, 16.91)$] varies as a function of energy. The cross section ratios for formation of $^8\text{Be}^*(17.64, 18.15)$ vary in a way consistent with a change in the population of the $T = 1$ part of the wave function over the energy range: at the higher energies, there is very little isospin violation. At higher $E_x$ only the $3^+$ state at $E_x = 19.2$ MeV is observed, the neighboring $3^+$ state at $E_x = 19.06$ MeV is not seen. The $J^\pi = 1^+$; $T = 0$ state is found to have $E_x = 18.146 \pm 0.005$ MeV (based on 17.638 for the $J^\pi = 1^+$; $T = 1$ state) and $\Gamma = 138 \pm 6$ keV (1970CA12). There is some question as to whether a two-level fit can be made for the $\alpha$ groups to $^8\text{Be}^*(16.63, 16.91)$. ((1970CA12 and W.D. Callender, private communication) are dubious about this, feeling that other $2^+$ states have to be brought into the calculation. Based on a two-level fit they find the following average values: $\Gamma_{16.6} = 113$ keV, $\Gamma_{16.9} = 80$ keV, $\Delta Q = 302$ keV. However, (1971IN004) state that the two-level fit is appropriate if the spectra are properly corrected for effects of final state Coulomb interactions: $\Gamma_{16.6} = 90 \pm 5$ keV, $\Gamma_{16.9} = 70 \pm 5$ keV, $\Delta Q = 290 \pm 7$ keV. See also (1966BR22) and (1970KI1D; theor.). For a listing of the parameters of observed states see Tables 8.4 and 8.5 (1961ER01, 1969NU01, 1970CA12).

Angular correlation studies [$E_d \leq 3$ MeV] indicate that reaction (b) takes place mainly by a sequential process involving $^8\text{Be}^*(0, 2.9, 11.4, 16.6, 16.9)$: see (1968LO01, 1970ST02, 1971LA14) and (1967CA13, 1968AS01). (1968LO01) report $E_x = 2.7 \pm 0.2$ MeV, $\Gamma = 1.0 \pm 0.1$ MeV, and $\Gamma = 3.0 \pm 0.5$ MeV for $^8\text{Be}^*(11.4)$. See also (1967CA13, 1970ST02). For a study of rescattering effects see (1972VA1L).


39. $^{10}\text{B}(^3\text{He}, p\alpha)^8\text{Be}$

At $E(^3\text{He}) = 2.45$ and 6.00 MeV this reaction proceeds primarily by sequential decay via $^8\text{Be}^*(0, 2.9)$ and via $^5\text{Li}$, $^9\text{B}$ and $^{12}\text{C}$ states [see also the latter nuclei] (1966WA16). See also (1966WI08, 1968KR02, 1970BE1F), (1966LA04, 1967HO1C) and (1967PR1B; theor.).

40. $^{10}\text{B}(\alpha, ^6\text{Li})^8\text{Be}$

At $E_\alpha = 46$ MeV angular distributions obtained for the transitions to $^8\text{Be}^*(0, 2.9)$ are consistent with a direct interaction mechanism (1970ZE03).

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41. (a) $^{11}\text{B}(\gamma, t)^{8}\text{Be}$  
\[ Q_{m} = -11.2242 \]
(b) $^{11}\text{B}(\gamma, t)^{4}\text{He}^{4}\text{He}$  
\[ Q_{m} = -11.1323 \]

See $^{11}\text{B}$ in (1975AJ02).

42. (a) $^{11}\text{B}(p, \alpha)^{8}\text{Be}$  
\[ Q_{m} = 8.591 \]
(b) $^{11}\text{B}(p, \alpha)^{4}\text{He}^{4}\text{He}$  
\[ Q_{m} = 8.6824 \]
\[ Q_{0} = 8.575 \pm 0.011 \] (1967SP09).

Angular distributions have been measured at $E_{p} = 0.78$ to 12.00 MeV (1963SY01, 1968WA1G; $\alpha_{0}$), 1.4, 2.0 and 2.6 MeV (1972GE19; $\alpha_{0}$, $\alpha_{1}$ (not at 2.0)), 12, 20, 24 and 30 MeV (1971CA16; $\alpha_{0}$), 26.7 and 38 MeV (1969GA03, 1970GU06; $\alpha_{0}$), 40 MeV (1971KA21; $\alpha_{0}$, $\alpha_{1}$ and $\alpha$ to $^{8}\text{Be}^{*}(12.5)$ [\(\Gamma = 4.0 \pm 0.5\) MeV] and to $^{8}\text{Be}^{*}(16.6 + 16.9, 17.6, 18.1, 19.0)$) and at $E_{p} = 45$ MeV (1971DE2B, 1972DE01, 1972DE02; $\alpha_{0}$, $\alpha_{1}$). At $E_{p} = 45$ MeV the angular distributions are typical of a direct reaction mechanism, with a rise in the backward direction indicative of heavy particle stripping (1972DE01, 1972DE02). Observed parameters for $^{8}\text{Be}^{*}(2.9)$ are shown in Table 8.4 (1969NU01, 1971KA21). At $E_{p} = 40$ MeV, $\theta = 20^\circ$, $d\sigma_{16.6}/d\sigma_{16.9} = 2.3 \pm 0.4$ (1971KA21).

Reaction (b) has been studied for $E_{p} = 0.15$ to 9.5 MeV. The reaction proceeds predominantly by sequential two-body decay via $^{8}\text{Be}^{*}(0, 2.9)$: see, e.g., (1965BR18, 1968CH01, 1972HU04). See also (1967KA09, 1967MA11, 1968GI03, 1969QU01, 1970CO03, 1971KO22, 1972MI1J). Some papers report very narrow widths in this reaction for $^{8}\text{Be}^{*}(2.9)$. However, (1972HU04) find a good fit to the data with $E_{x} = 2.99$ MeV, $\Gamma = 1.45$ MeV when an interference term is included. The interference effect is attributable to the identity of the three $\alpha$-particles and to the ambiguity in their order of emission (1965BR18).


43. $^{11}\text{B}(d, n\alpha)^{8}\text{Be}$  
\[ Q_{m} = 6.3658 \]

See (1971RE19) and $^{9}\text{Be}$.

44. $^{11}\text{B}(^{3}\text{He}, ^{6}\text{Li})^{8}\text{Be}$  
\[ Q_{m} = 4.570 \]

This reaction has been studied for $E(^{3}\text{He}) = 1.4$ to 5.8 MeV. Angular distributions have been measured at $E(^{3}\text{He}) = 5.2$ MeV involving $^{8}\text{Be}_{g.s.} + ^{6}\text{Li}_{g.s.}$, $^{8}\text{Be}_{g.s.} + ^{6}\text{Li}_{3.56}^{*}$, and $^{8}\text{Be}_{2.9}^{*} + ^{6}\text{Li}_{g.s.}$ (1967YO02, 1967YO1C, 1968ME13).
45. (a) \( ^{11}\text{B}(\alpha, ^7\text{Li})^8\text{Be} \) \( Q_m = -8.758 \)
(b) \( ^{11}\text{B}(\alpha, ^7\text{Li})^4\text{He}^4\text{He} \) \( Q_m = -8.666 \)

Angular distributions have been reported at \( E_\alpha = 28.4 \) and \( 29.0 \) MeV for \( ^8\text{Be}_{g.s.} + ^7\text{Li}_{g.s.} \), \( ^8\text{Be}_{g.s.} + ^7\text{Li}_{*_{0.48}} \) and \( ^8\text{Be}_{*_{2.9}} + ^7\text{Li} \) (29 MeV only) by (1968KA24) and at 42 MeV for \( ^8\text{Be}_{g.s.} + ^7\text{Li}_{g.s.} \) and \( ^8\text{Be}_{g.s.} + ^7\text{Li}_{*_{0.48}} \) by (1968MI05). At \( E_\alpha = 65 \) MeV \( ^8\text{Be}^*_{(16.6 + 16.9, 20.0)} \) are apparently also excited (1973WO06). See also (1966GE12) and \( ^7\text{Li} \). For reaction (b) see (1969FU09).

46. (a) \( ^{12}\text{C}(\gamma, \alpha)^8\text{Be} \) \( Q_m = -7.367 \)
(b) \( ^{12}\text{C}(e, e\alpha)^8\text{Be} \) \( Q_m = -7.367 \)

For reaction (a) see (1966LA04), (1973CL1E) and (1965DZ1A; theor.). For reaction (b) see (1970EN1A).

47. \( ^{12}\text{C}(n, n\alpha)^8\text{Be} \) \( Q_m = -7.367 \)

This reaction proceeds via \( ^8\text{Be}^*_{(0, 2.9)} \) at \( E_n = 13 \) to \( 18 \) MeV, and via states in \( ^5\text{He}, ^9\text{Be} \) and \( ^{12}\text{C} \) (1966MO05). See also (1966LA04) and (1971DO1K).

48. \( ^{12}\text{C}(p, p\alpha)^8\text{Be} \) \( Q_m = -7.367 \)

This reaction has been studied for \( 13 \leq E_p \leq 160 \) MeV. At low energies it involves \( ^8\text{Be}(0) \); at higher energies \( ^8\text{Be}^*_{(0, 2.9)} \) (1966RO1D, 1967GA01, 1969LU1B, 1970GO12, 1970KE1B, 1972MA62). It is not clear whether higher states are also involved: see (1970KE1B). See also (1966LA04) and (1966JA1B, 1968YA1C, 1972YA1B).

49. \( ^{12}\text{C}(d, ^6\text{Li})^8\text{Be} \) \( Q_m = -5.893 \)

Angular distributions have been determined at \( E_d = 19.5 \) MeV (1971GU07), 28 MeV (1972BE1T, 1972BE29: \( ^8\text{Be}(0) \)) and 51.8 MeV (1970EI05; both \( ^8\text{Be}^*_{(0, 2.9)} \)). At \( E_d = 28 \) MeV a structure is observed which is attributed to the process \( ^{12}\text{C}(d, \alpha)^{10}\text{B} \rightarrow \alpha + ^6\text{Li} \) (1972CO23). At \( E_d = 55 \) MeV the population of \( ^8\text{Be}^*_{(11.4, 16.6 + 16.9)} \) is also reported (1971MC04). See also (1966DA1C, 1970AN1E), (1967OG1A, 1972GA1E) and (1968RO1D, 1970EL1F, 1971DR02, 1973HE1J; theor.). See also \( ^{12}\text{C} \) and (1966LA04).
50. $^{12}$C($^3$He, $^7$Be)$^8$Be  \[ Q_m = -5.780 \]

Angular distributions have been measured for the transitions to $^7$Be$_{g.s.} + ^8$Be$_{g.s.}$ and $^7$Be$^*_0.43 + ^8$Be$_{g.s.}$ at $E(^3$He) = 25.5 to 29 MeV ($1972$PI$1A$, $1973$PI$1B$, $1973$PI$1D$), 28 MeV ($1970$DE$12$, $1973$KL$1B$) and at 35.7 MeV ($1969$ZE$1A$, $1970$FO$1D$). The transitions to $^7$Be$_{g.s.} + 0.43 + ^8$Be$^*_2.9$ have also been studied by ($1970$DE$12$). See also ($1967$ZA$1B$, $1973$ST$1N$) and ($1969$NE$1D$; theor.).

51. (a) $^{12}$C($\alpha$, 2$\alpha$)$^8$Be  \[ Q_m = -7.367 \]
(b) $^{12}$C($\alpha$, $^8$Be)$^8$Be  \[ Q_m = -7.4587 \]

This reaction has been studied up to $E_\alpha = 104$ MeV. At $E_\alpha = 25$ MeV it involves $^8$Be(0) ($1966$BO$28$); at $E_\alpha = 28.0, 37.4, 70, 90$ and 104 MeV, the reaction goes via $^8$Be$^*(0, 2.9)$ ($1965$YA$02$, $1967$TA$1C$, $1968$YA$02$, $1971$BR$1G$, $1972$SH$1J$) and at 90 MeV it may, in addition, involve the broad 4$^+$ state at 11.4 MeV ($1970$JA$06$). See also ($1966$LA$04$) and ($1965$KU$1B$, $1967$ME$1C$; theor.).

Reaction (b) has been studied for $11.9 \leq E_\alpha \leq 19.4$ MeV and angular distributions are reported for $E_\alpha = 12.70$ to 16.25 MeV ($1967$CH$21$) and 65 MeV ($1973$WO$06$: $^8$Be$^*(0, 2.9)$). See $^{16}$O in ($1971$AJ$02$).

52. $^{12}$C($^{12}$C, $^{16}$O)$^8$Be  \[ Q_m = -0.2051 \]

Angular distributions have been measured for $E(^{12}$C) = 11.6 to 13.4 MeV ($1972$CO$1H$). At $E(^{12}$C) = 50 to 65 MeV the population of $^8$Be$^*(0, 2.9)$ is reported by ($1972$FL$1C$). See also ($1968$JA$1F$, $1970$JA$1B$, $1972$GR$1T$, $1973$CR$1A$, $1973$SC$1J$).

53. $^{12}$C($^{14}$N, $^{18}$F)$^8$Be  \[ Q_m = -2.951 \]

See ($1965$WI$1A$).

54. $^{12}$C($^{16}$O, $^{20}$Ne)$^8$Be  \[ Q_m = -2.637 \]

See ($1972$GR$1Q$).

55. $^{13}$C(p, $^6$Li)$^8$Be  \[ Q_m = -8.615 \]
Angular distributions have been measured at $E_p = 45$ MeV for the transitions to $^8\text{Be}^*(0, 2.9)$ (1971BR07).

56. $^{13}\text{C}(d, ^7\text{Li})^8\text{Be}$ \hspace{1cm} $Q_m = -3.589$

Angular distributions are reported at $E_d = 14.6$ MeV for the transitions to $^8\text{Be}_{g.s.} + ^7\text{Li}_{g.s.}$ and $^8\text{Be}_{g.s.} + ^7\text{Li}^*_{0.48}$ (1967DE03).

57. $^{13}\text{C}(^3\text{He}, ^8\text{Be})^8\text{Be}$ \hspace{1cm} $Q_m = 8.173$

Angular distributions have been obtained at $E(\text{He}) = 3.3, 5.0$ and 5.8 MeV for the transition to $^8\text{Be}_{g.s.} + ^8\text{Be}_{g.s.}$ (1968JA07, 1969JA1L). See also $^{16}\text{O}$ in (1971AJ02) and (1967SA1E).

58. (a) $^{14}\text{N}(n, ^7\text{Li})^4\text{He}^4\text{He}$ \hspace{1cm} $Q_m = -8.823$
(b) $^{14}\text{N}(n, t)^4\text{He}^4\text{He}^4\text{He}$ \hspace{1cm} $Q_m = -11.290$

See (1967MO21, 1971SC16) and $^7\text{Li}$.

59. $^{14}\text{N}(^{10}\text{B}, 6\alpha)$ \hspace{1cm} $Q_m = 0.366$


60. $^{16}\text{O}(\gamma, 2\alpha)^8\text{Be}$ \hspace{1cm} $Q_m = -14.528$

See (1966LA04) and $^{12}\text{C}$ in (1968AJ02).

61. $^{16}\text{O}(p, p)^4\text{He}^4\text{He}^4\text{He}^4\text{He}$ \hspace{1cm} $Q_m = -14.436$

See (1961KO02, 1962VA1A) and $^{16}\text{O}$ in (1971AJ02).

62. $^{16}\text{O}(\alpha, ^{12}\text{C})^8\text{Be}$ \hspace{1cm} $Q_m = -7.2535$
Angular distributions have been measured at $E_\alpha = 65$ MeV involving $^8\text{Be}_{g.s.}$ and $^{12}\text{C}^*(0, 4.4)$ (1973WO06). See also (1968PA12) and $^{16}\text{O}$ in (1971AJ02) and $^{20}\text{Ne}$ in (1972AJ02).

63. $^{16}\text{O}(^{10}\text{B}, 6\alpha)$  $Q_m = 7.234$

See (1965SH11).

64. $^{19}\text{F}(p, ^{12}\text{C})^8\text{Be}$  $Q_m = 0.861$

See (1969GO1B, 1971GO1U) and $^{20}\text{Ne}$ in (1972AJ02).

65. $^{20}\text{Ne}(\alpha, ^{16}\text{O})^8\text{Be}$  $Q_m = -4.822$

This reaction has not been observed: see (1962LA15).
8B
(Figs. 13 and 14)

GENERAL: (See also (1966LA04).)


Special levels: (1966BA26).


\[ \mu = 1.0355 \pm 0.003 \text{ nm} \ (1973MIYZ). \]

1. \(^{8}\text{B}(\beta^{+})^{8}\text{Be} \quad Q_{m} = 17.981\]

The \(\beta^{+}\) decay leads mainly to \(^{8}\text{Be}^{*}(2.9)\). The half-life is \(774 \pm 5 \text{ msec} \ (1964MA35)\), \(762 \pm 5 \text{ msec} \ (1971WI05)\), \(772 \pm 4 \text{ msec} \ (1973MCZW)\): the mean is \(770 \pm 3 \text{ msec} \). See also Table 8.19 in (1966LA04): \(\log ft = 5.64 \ (1966BA1A)\) \(^{\dagger}\). There is also a branch to a \(^{8}\text{Be}\) state at \(\approx 16.6 \text{ MeV; } \log ft = 2.9 \ (1964MA35)\). See, however, (1969BA43): \(\log ft = 3.33\). Measurements of the excitation spectra in the decays of \(^{8}\text{Li}\) and \(^{8}\text{B}\) show no evidence for second-class currents: \(|g_{IT}| < 7 \times 10^{-4} \ (1971WI05)\). See also (1964FO1A, 1969BA1X, 1969BA1M, 1972BA2M, 1972KO1A, 1973BA2C) for astrophysical applications, (1967BA1M, 1970DA21, 1971LI1H, 1971WI1C, 1971WI18, 1972EM02, 1972HO23, 1972WI28, 1972WI1C, 1973EM1B, 1973HA49, 1973TO14, 1974WI1L; theor.).

2. \(^{6}\text{Li}(^{3}\text{He}, n)^{8}\text{B} \quad Q_{m} = -1.975\]

At \(E(^{3}\text{He}) = 3.5\) to \(5.7 \text{ MeV}\), time-of-flight spectra locate the first excited state at \(0.767 \pm 0.012 \ (1965FA03, 1966FA1A)\), \(0.783 \pm 0.010 \text{ MeV} \ (1965MC06; \Gamma = 40 \pm 10 \text{ keV})\). Angular distributions for the \(n_{0}\) group have been reported at \(E(^{3}\text{He}) = 4.8\) to \(5.7 \text{ MeV}\). The appearance of a forward peak indicates an \(L = 0\) transfer and hence a knockout mechanism (\(L = 0\) is forbidden for simple diproton stripping) (1967VA24). See also (1963DI02). A state at \(E_{x} = 2.17 \pm 0.05 \text{ MeV}\) is reported by (1963DI02; abstract). For threshold measurements see (1966LA04). See also \(^{9}\text{B}\).

\(^{\dagger}\) and B. Zimmerman, private communication.
Table 8.11: Energy levels of $^8$B

<table>
<thead>
<tr>
<th>$E_x$ (MeV ± keV)</th>
<th>$J^\pi; T$</th>
<th>$\tau_{1/2}$ or $\Gamma_{\text{c.m.}}$ (keV)</th>
<th>Decay</th>
<th>Reactions</th>
</tr>
</thead>
<tbody>
<tr>
<td>g.s.</td>
<td>2$^+$; 1</td>
<td>$\tau_{1/2} = 769 \pm 4$ msec</td>
<td>$\beta^+$</td>
<td>1, 2, 3, 4, 5, 6, 7</td>
</tr>
<tr>
<td>0.778 ± 7</td>
<td>3$^+$; 1</td>
<td>$\Gamma = 40 \pm 10$</td>
<td>$\gamma, p$</td>
<td>2, 3, 4, 5</td>
</tr>
<tr>
<td>2.32 ± 30</td>
<td></td>
<td>350 ± 40</td>
<td></td>
<td>2, 4, 5</td>
</tr>
</tbody>
</table>

3. $^7$Be(p, $\gamma$)$^8$B  

Absolute cross sections have been measured for $E_p = 0.165$ to 10.0 MeV (1969KA1K, 1972KA1B), 0.48 to 1.93 MeV (1966PA16) and 0.95 to 3.28 MeV (1970VA26). In the data of (1966PA16, 1972KA1B) a resonance at $E_p = 724$ keV reflects $^8$B*(0.77) with $\Gamma_{\text{lab}} \approx 46$ keV: the peak cross section is $2.20 \pm 0.22 \mu$b, $\Gamma_\gamma = 50 \pm 25$ meV. See also (1969KA1K).

The low-energy cross-section factor, evaluated at the Gamow peak ($E_p \approx 20$ keV), $S(0.02) = 0.031$ keV · b (1973RO08). Other values are $S(0) = 0.0335 \pm 0.003$ keV · b, $dS/dE = -3 \times 10^{-5}$ b (1972KA1B), $0.035 \pm 0.004$ keV · b, $dS/dE = -3 \times 10^{-5}$ b (1968PA1M) and $0.0430 \leq S(0) \leq 0.0453$ keV · b (1970AU1B). See also (1969KA1K, 1970VA26). The relevance of this reaction to astrophysics is discussed by (1965TO02, 1966PA16, 1967FO1B, 1967TO1B, 1968BA1W, 1968BA2E, 1968PA1M, 1969BA1U, 1969BA1M, 1970AU1B, 1970VA26, 1971BA2X, 1972BA2M, 1972KA1B, 1972PA1C, 1972TO1D, 1973BA2C, 1973RO08, 1973TR1E).

4. $^{10}$B(p, t)$^8$B  

$Q_m = -18.531$

At $E_p = 49.5$ MeV angular distributions have been measured for the tritons to $^8$B*(0, 2.32) (1970SQ01): $L = 2$ and $L = 0 + 2$, leading to $J^\pi = 2^+$ and $3^+$, respectively. See also $^{10}$B(p, $^3$He)$^6$Be. The energy of the excited state is $2.29 \pm 0.05$ MeV (1970SQ01), $2.34 \pm 0.04$ MeV (1968BR23): $\Gamma_{\text{lab}} = 390 \pm 40$ keV (1967MC14). $^8$B*(0.78) is also observed. See also (1971KA04; theor.).

5. $^{11}$B($^3$He, $^6$He)$^8$B  

$Q_m = -16.920$

At $E(^3$He) = 50 MeV, $^6$He groups are observed to the first three states of $^8$B (1967MC14).

6. $^{12}$C($^3$He, $^7$Li)$^8$B  

$Q_m = -22.899$

This reaction has been studied at $E(^3$He) = 40.7 MeV (1971DE37).
7. $^{12}\text{C}(\alpha, {}^8\text{Li})^8\text{B}$  \hspace{1cm} $Q_m = -41.445$

See (1968MC02).

$^8\text{C}$

(Not illustrated)

$^8\text{C}$ has been observed in the $^{12}\text{C}(\alpha, {}^8\text{He})^8\text{C}$ reaction at $E_\alpha = 156$ MeV; $M - A = 35.30 \pm 0.20$ MeV, $\Gamma_{\text{c.m.}} = 220^{+80}_{-140}$ keV [the differential cross section at $2^\circ$ (lab) is $\approx 20 \text{ nb/sr}$] (R.G.H. Robertson, S. Martin, W.R. Falk, D. Ingham and A. Djalois, private communication). $^8\text{C}$ is then unstable with respect to $^7\text{B} + p (Q = 0.1), ^6\text{Be} + 2p (Q = 2.3), ^5\text{Li} + 3p (Q = 1.8), ^4\text{He} + 4p (Q = 3.7)$. See also (1960GO1B, 1966KE16, 1970WA1G).
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(Closed December 31, 1973)

References are arranged and designated by the year of publication followed by the first two letters of the first-mentioned author’s name and then by two additional characters. Most of the references appear in the National Nuclear Data Center files (Nuclear Science References Database) and have NNDC key numbers. Otherwise, TUNL key numbers were assigned with the last two characters of the form 1A, 1B, etc. In response to many requests for more informative citations, we have, when possible, included up to ten authors per paper and added the authors’ initials.

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